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Magnetoelastic coupling and Grüneisen scaling in NdB₄

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We report high-resolution capacitance dilatometry studies on the uniaxial length changes in a NdB₄ single crystal. The evolution of magnetically ordered phases below $T_N = 17.2$ K [commensurate antiferromagnetic phase (cAFM)], $T_{IT} = 6.8$ K [intermediate incommensurate phase (IT)], and $T_{LT} = 4.8$ K [low-temperature phase (LT)] is associated with pronounced anomalies in the thermal expansion coefficients. The data imply significant magnetoelastic coupling and evidence of a structural phase transition at T_{LT} . While both cAFM and LT favor structural anisotropy δ between in-plane and out-of-plane length changes, it competes with the IT type of order, i.e., δ is suppressed in that phase. Notably, finite anisotropy well above T_N indicates short-range correlations which are, however, of neither cAFM, IT, nor LT type. Grüneisen analysis of the ratio of thermal expansion coefficient and specific heat enables the derivation of uniaxial as well as hydrostatic pressure dependencies. While α/c_p evidences a single dominant energy scale in LT, our data imply precursory fluctuations of a competing phase in IT and cAFM, respectively. Our results suggest the presence of orbital degrees of freedom competing with cAFM, and successive evolution of a magnetically and orbitally ordered ground state.

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I. INTRODUCTION

Geometric frustration provides a prime route to macroscopic degeneracy of ground states, thereby suppressing or even completely eliminating static long-range order [1]. Often, when magnetic order eventually evolves, competing interactions result in complex magnetic phase diagrams constituting energetically similar spin configurations. Prototypical examples of geometrically frustrated magnets are triangular, kagome, and pyrochlore lattices [2–10]. While many of these systems realize quantum magnets, rare earth (R) borides, due to the large magnetic moment of the rare earth ions, are well suited for investigating frustrated magnets in a classical limit. The family of tetraborides RB₄ is one particular example exhibiting magnetic frustration due to the rare earth ions being arranged on the geometrically frustrated Shastry-Sutherland lattice (SSL) [11,12]. Many of the RB₄ compounds display a variety of competing phases at low magnetic fields [12]. As with all RB₄ compounds featuring trivalent rare earth ions, NdB₄ reported here is metallic [13]. As such the Ruderman-Kittel-Kasuya-Yosida (RKKY) interaction may play an important role in addition to competing interactions evolving on exchange paths formed by the SSL.

The RB₄ family crystallizes in a tetragonal D_{4h}⁵-P4/mbm structure [14]. The boron sublattice is made up of octahedra forming chains along the crystallographic *c* direction. These

chains are connected via two boron atoms to form rings within the *ab* plane. The rare earth ions are located above and below the middle of these rings, constituting the rare earth sublattice. In NdB₄, magnetism is due to Nd³⁺ ions whose electron configuration $4f^3$ implies a $^4I_{9/2}$ ground state. Upon cooling, NdB₄ shows three successive phase transitions, signaling the onset of commensurate antiferromagnetic (cAFM) order at $T_N = 17.2$ K, an incommensurate antiferromagnetic phase (IT) evolving at $T_{IT} = 6.8$ K and a low-temperature phase (LT) below $T_{LT} = 4.8$ K [15,16]. The former phase features an all-in structure of magnetic moments pointing into the squares formed by the Nd ions [17]. In addition, the moments are slightly tilted towards the *c* axis [17]. The pseudoquartet ground state consisting of two Kramers doublets carries an electric quadrupole moment, giving rise to orbital degrees of freedom [18]. Here, we study the effects of long-range order on the length changes in NdB₄ by high-resolution capacitance dilatometry. Despite the relevance of dilatometry for elucidating the coupling of lattice, orbital, and magnetic degrees of freedom, particularly in rare-earth systems, such studies are currently missing for the whole family of tetraborides RB₄. We find pronounced thermal expansion anomalies at the magnetic transitions as well as evidence for a structural transition at T_{LT} . In addition, precursory fluctuations well above T_N associated with negative in-plane thermal expansion indicate short-range ordering in the paramagnetic phase which is, however, not associated with an ordering phenomena evolving in either of the three low-temperature phases. Analysis of the Grüneisen ratios and of structural anisotropy enables elucidating the interplay of lattice, structure, and presumably orbital degrees of freedom.

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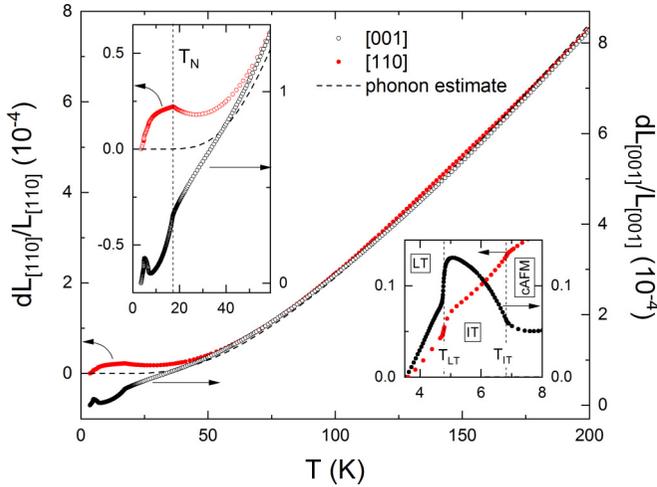


FIG. 1. Relative length changes dL_i/L_i along the crystallographic [001]- (black circles) and [110]-direction (red circles) vs. temperature. Insets show enlargement of the low temperature region. Left and right ordinates are to scale but shifted in order to account for the difference at 200 K. Vertical dashed lines indicate the transition temperatures T_N , T_{IT} , and T_{LT} toward the commensurate antiferromagnetic phase (cAFM), the incommensurate IT phase, and the low-temperature phase LT.

II. EXPERIMENTAL METHODS

Single crystals of NdB_4 were grown by the optical floating-zone technique as reported in detail in Ref. [19]. The relative length changes dL_i/L_i along the crystallographic [001] and [110] directions (space group 127), respectively, were studied on an oriented cuboid-shaped single crystal of dimensions $1.476 \times 1.478 \times 1.880 \text{ mm}^3$. The measurements were done by means of a three-terminal high-resolution capacitance dilatometer [20], and the linear thermal expansion coefficients $\alpha_i = 1/L_i \times dL_i(T)/dT$ were derived. Static magnetic susceptibility $\chi = M/B$ was studied in a magnetic field of $B = 0.1 \text{ T}$ in a vibrating sample magnetometer (VSM) of a Magnetic Properties Measurement System (MPMS3 SQUID magnetometer) by Quantum Design. Specific heat measurements have been performed in a Quantum Design PPMS using a relaxation method.

III. LENGTH CHANGES UPON EVOLUTION OF MAGNETIC ORDER

Thermal expansion in NdB_4 is positive at high temperatures and shows similar behavior in-plane and out-of-plane (see Fig. 1). Upon cooling below $\simeq 50 \text{ K}$, $L_{[001]}$ decreases significantly stronger compared to $L_{[110]}$. Distinct anomalies in the thermal expansion at T_N as well as at T_{LT} and T_{IT} imply pronounced magnetoelastic coupling.

Magnetoelastic effects associated with the various ordering phenomena are particularly evident if the thermal expansion coefficients are considered, which display pronounced anomalies at $T_N = 17.2 \text{ K}$, $T_{IT} = 6.8 \text{ K}$, and $T_{LT} = 4.8 \text{ K}$ [see Figs. 2 and 3(a)]. The onset of long-range antiferromagnetic order at T_N is confirmed by previous magnetization and neutron diffraction data [15–17], while two distinct antiferromagnetic

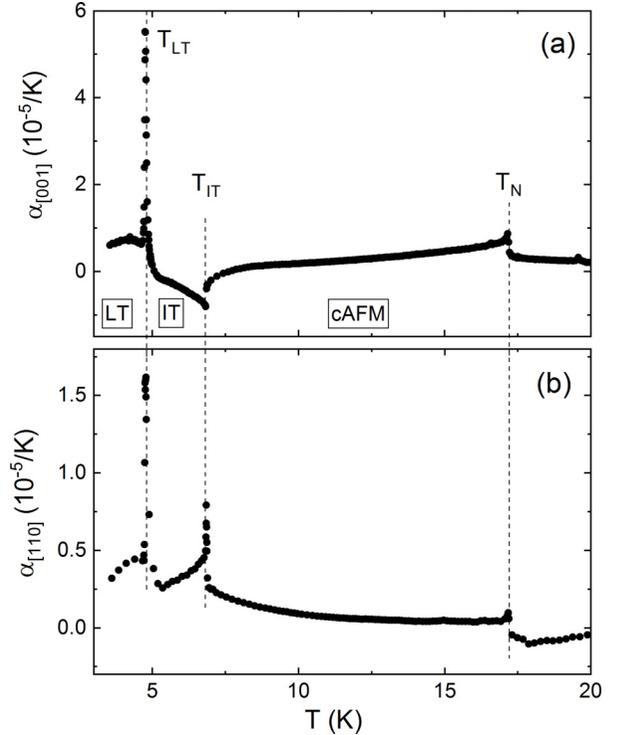


FIG. 2. Uniaxial thermal expansion coefficients α_i along the crystallographic (a) [001] and (b) [110] direction. Dashed lines indicate the transition temperatures T_{LT} , T_{IT} , and T_N (see the text).

phases develop at T_{LT} and T_{IT} [17]. At T_N and T_{IT} , the thermal expansion data imply λ -shaped anomalies typical for continuous phase transitions. Anomalies in the volume thermal expansion coefficient $\alpha_v = 2\alpha_{\perp c} + \alpha_{\parallel c}$ show similar features [Fig. 3(a)]. All anomalies at T_N imply the decrease in length upon onset of magnetic order, i.e., there is positive uniaxial pressure dependence $dT_N/dp_i > 0$ for both $i = [001]$ and [110]. Accordingly, hydrostatic pressure dependence of T_N is positive too.

In contrast, anomalies in the uniaxial thermal expansion coefficients exhibit opposite signs at T_{IT} , i.e., $dT_{IT}/dp_{\parallel c} < 0$ and $dT_{IT}/dp_{\perp c} > 0$ (Fig. 2). The resulting volume response [Fig. 3(a)] exhibits a positive sign of anomaly and hence positive hydrostatic pressure dependence. At T_{LT} there are jumps ΔL_i in the observed length changes for both directions i , which imply discontinuous evolution of the LT phase. The corresponding jumps in $L_i(T)$ are negative upon cooling (see the insets of Fig. 1), which implies $dT_{LT}/dp_i > 0$ in both cases. Quantitatively, at T_{LT} we observe $\Delta L_{[110]} = -2.0(2) \times 10^{-6}$ and $\Delta L_{[001]} = -5.0(2) \times 10^{-6}$.

In order to quantitatively evaluate the effect of magnetic order and crystal field on thermal expansion, phonon contributions are estimated from specific heat data on LaB_4 [21] by assuming Grüneisen scaling of the phonon heat capacity. Specifically, the phonon specific heat of LaB_4 is given in [21] as

$$c_p^{ph} = a_1 D\left(\frac{T}{\Theta_{D1}}\right) + a_2 D\left(\frac{T}{\Theta_{D2}}\right) + a_3 E\left(\frac{T}{\Theta_{E1}}\right) + a_4 E\left(\frac{T}{\Theta_{E2}}\right). \quad (1)$$

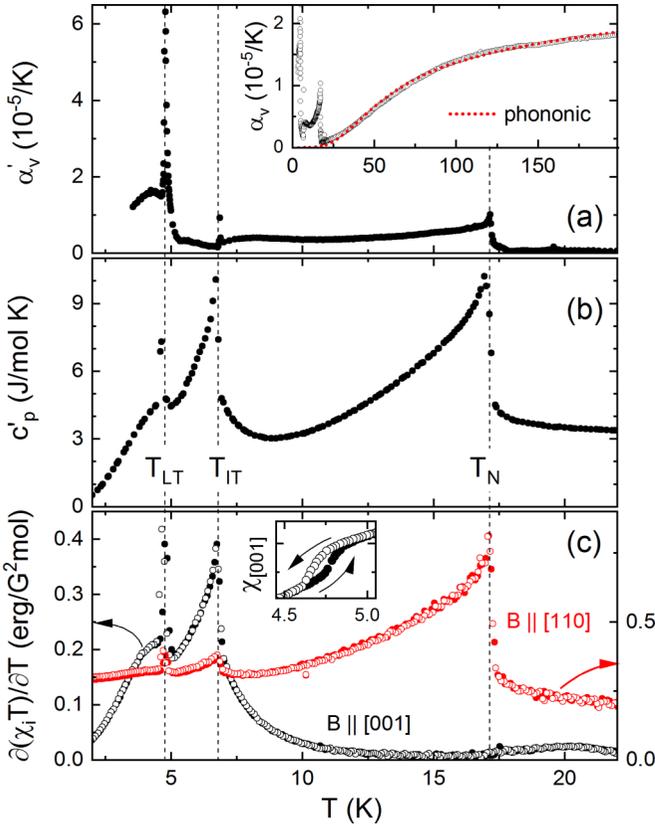


FIG. 3. (a) Volumetric thermal expansion coefficient corrected by the phonon contribution. Inset: Uncorrected volumetric thermal expansion coefficient and phonon fit to the high temperature data (see the text). (b) Corresponding background corrected specific heat, and (c) Fisher's specific heat $\partial(\chi_i T)/\partial T$ for $B = 0.1$ T applied along the crystallographic [100] (black circles) and [110] direction (red circles), respectively. Filled (open) circles indicate measurements upon heating (cooling). Inset: Corresponding magnetic susceptibility $\chi_{[001]}$ highlighting hysteresis at T_{LT} . Vertical dashed lines indicate T_{LT} , T_{IT} , and T_N .

Here, a_i are parameters, and D and E denote Debye and Einstein functions with Debye and Einstein temperatures as given in Table I. The same function was fitted to the uniaxial thermal expansion coefficients $\alpha_{\parallel c}$ and $\alpha_{\perp c}$ above 65 K, using the same Einstein temperatures and allowing the parameters a_i to vary. The Debye temperatures were multiplied by the correction factor $r=0.9828$ to account for the different masses

TABLE I. Parameters of phonon background of thermal expansion. Θ_D and Θ_E are Debye and Einstein temperatures, describing a phonon background of specific heat in LaB_4 from Ref. [21]. The parameters γ are the uniaxial and volume phononic Grüneisen parameters.

γ (mol/MJ)			
$\Theta_{D1} = 420$ K	$\gamma_{D1}^{\parallel c} = 0.055$	$\gamma_{D1}^{\perp c} = 0.069$	$\gamma_{D1}^v = 0.191$
$\Theta_{D2} = 230$ K	$\gamma_{D2}^{\parallel c} = 2.27$	$\gamma_{D2}^{\perp c} = 0$	$\gamma_{D2}^v = 2.27$
$\Theta_{E1} = 890$ K	$\gamma_{E1}^{\parallel c} = 0.020$	$\gamma_{E1}^{\perp c} = 0.013$	$\gamma_{E1}^v = 0.066$
$\Theta_{E2} = 177$ K	$\gamma_{E2}^{\parallel c} = 0$	$\gamma_{E2}^{\perp c} = 0.349$	$\gamma_{E2}^v = 0.682$

of Lanthanum and Neodymium [22]. The volumetric thermal expansion coefficient obtained by $\alpha_v = 2\alpha_{\perp c} + \alpha_{\parallel c}$ was fitted in the same way, except that a_2 was set to the value obtained from the fit to $\alpha_{\parallel c}$. From these fits the phononic Grüneisen parameters $\gamma = \alpha/c_p$ are determined for both uniaxial thermal expansion coefficients and for the volume as listed in Table I.

The result of this estimate is shown by the dashed line in Fig. 1. It implies anisotropic nonphononic length changes up to about 65 K, while for higher temperatures the experimental data are well described by the fits (see Fig. 1).

Figure 3(a) shows the volumetric thermal expansion coefficient α'_v corrected by the phonon contribution. The inset shows the bare experimental data α_v and the phonon fit derived as described above. In the whole temperature regime under study, α'_v is positive and in particular, as described above, all anomalies are positive indicating positive hydrostatic pressure dependencies $dT_i/dp > 0$ of all transition temperatures. In addition, Fig. 3(b) displays the background corrected specific heat c'_p obtained by analogously subtracting the phonon entropy changes. For comparison, we also present Fisher's specific heat [23] $c_p^m \propto \partial(\chi_i T)/\partial T$ for magnetic fields applied $B \parallel c$ and $B \perp c$ [Fig. 3(c)], which provides an estimate of magnetic entropy changes. Analogously to α , c'_p and c_p^m display λ -shaped and discontinuous anomalies at the respective ordering temperatures. Note that in general, c_p^m resembles c'_p very well, which strongly confirms the procedure applied for background correction. The fact that there is no visible anomaly in $\partial(\chi_c T)/\partial T$ at T_N implies the easy axis being $\perp c$ in the cAFM phase which agrees to recent neutron data [17,18]. In addition, the magnetic entropy changes are well detected by $\chi_{\parallel c}$ in the IT and LT phases, suggesting significant out-of-plane spin components. We also note a small but finite hysteresis between cooling and heating in the magnetic susceptibility associated with the discontinuity at T_{LT} [see the inset of Fig. 3(c)].

IV. DISCUSSION

Comparing the anomalous length and entropy changes provides quantitative information on coupling between spin and structure. In the presence of one dominant energy scale ϵ , the ratio of the thermal expansion coefficient and specific heat is independent on temperature and enables the determination of the pressure dependence of ϵ by evaluating the volume or uniaxial Grüneisen parameter [24,25]:

$$\gamma_i = \frac{\alpha_i}{c_p} = \frac{1}{V_m} \left. \frac{\partial \ln \epsilon}{\partial p_i} \right|_T. \quad (2)$$

Here, V_m is the molar volume and the index i indicates a uniaxial direction or the volume. Exploiting the nonphononic contributions $\alpha'_{[110]}$, $\alpha'_{[001]}$, α'_v , and c'_p yields the Grüneisen parameters in Fig. 4. The experimentally obtained Grüneisen ratios distinctly change at the different phase boundaries, thereby signaling clear changes of magnetoelastic coupling. At T_N we observe jumps in γ_v and $\gamma_{\perp c}$, but no clear anomaly appears in γ_c . This absence of a clear anomaly corresponds to the absence of an anomaly in the magnetic susceptibility $\chi_{[001]}$ at T_N . Within the cAFM phase, both the uniaxial Grüneisen parameters $\gamma_{\parallel c}$ and $\gamma_{\perp c}$, as well as the volume Grüneisen parameter γ_v , are rather constant but begin to change a few

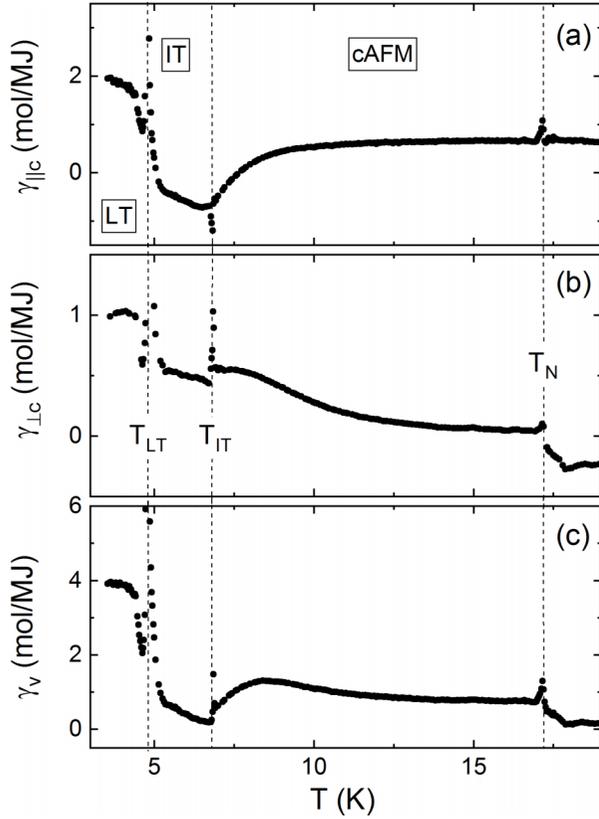


FIG. 4. Grüneisen parameters $\gamma_i = \alpha'_i/c'_p$, where i denotes the two directions (a) [001] and (b) [110], as well as the volume (c). Dashed lines show the transition temperatures T_N , T_{IT} , and T_{LT} .

Kelvin above T_{IT} . We conclude that critical fluctuations start to evolve at around 10 K, associated with structural changes of opposite uniaxial pressure dependence. The presence of opposite signs of anomalies tends toward volume conservation at the phase transition and is also a feature of the thermal expansion anomalies at T_{IT} (see Fig. 2). We hence conclude that fluctuations below 10 K are of the IT type. Notably, the changes of Grüneisen scaling in the cAFM phase correspond to the evolution of the purely magnetic (1 0 0) reflection, which gradually evolves below around 10 K and sharply increases at T_{LT} [17]. Microscopically, the fluctuation regime in cAFM seems to be associated with the rotation of magnetic

moments out of the ab plane as the out-of-plane angle is found to evolve below ca. 10 K as a secondary order parameter. In the IT phase, all parameters γ_i increase upon cooling. We again attribute this to precursing fluctuations — here of the LT type as corroborated by the equal signs of associated pressure dependencies — superimposed to the response of the IT type of order. Lastly there is a (non-volume-conserving) jump at the transition into the LT phase which features constant values of γ_i .

Our analysis yields the γ values summarized in Table II. Using the Ehrenfest relation for the continuous phase transitions at T_N and T_{IT} , the obtained values of γ yield the pressure dependencies of the ordering temperatures at vanishing pressure, i.e., $dT_{N/IT}/dp_i = T_{N/IT}V_m\gamma_i$ [25]. Here the pressure dependence of T_{IT} is estimated by the value of the Grüneisen parameter directly below the phase transition. In an attempt to obtain the pressure dependencies of T_N and T_{IT} directly from the anomalies at the phase transitions via the Ehrenfest relation $dT_{N/IT}/dp_i = T_{N/IT}V_m\Delta\alpha/\Delta c_p$, we deduced the actual jumps in the thermal expansion coefficients $\Delta\alpha$ and in the specific heat Δc_p at both transitions by means of area conserving constructions. While $\Delta\alpha_{IT,V}$ could not be extracted from the data, for the other directions the procedure yields similar pressure dependencies as obtained by exploiting the Grüneisen parameters (see Table II). For the discontinuous phase transition at T_{LT} , the pressure dependence is governed by the respective discontinuities in length, $\Delta L_{i,LT}/L_i$ and in entropy, ΔS_{LT} , via the Clausius-Clapeyron equation:

$$\frac{dT_{LT}}{dp_i} = V_m \frac{\Delta L_{i,LT}/L_i}{\Delta S_{LT}}, \quad (3)$$

with the molar volume $V_m = 3.22 \times 10^{-5} \text{ m}^3$ [19]. Analysis of the experimental specific heat data yields $\Delta S_{LT} = 0.13(3) \text{ J}/(\text{mol K})$. We note a general experimental uncertainty due to the relaxation method used for determining the heat capacity at a discontinuous phase transition. From the field dependence $dT_{LT}/dB \approx 0.4(1) \text{ K}/\text{T}$ obtained around 1 T and the jump in magnetization $\Delta M_{LT} \approx 85(10) \text{ erg}/(\text{mol} \times \text{G})$ we obtain $\Delta S_{LT} = (dT_{LT}/dB)^{-1} \times \Delta M_{LT} \approx 0.21(8) \text{ J}/(\text{mol K})$, which is hence used for the Clausius-Clapeyron analysis. The results for all phase transitions are shown in Table II.

The resulting pressure dependencies are small and the hydrostatic effects on T_N and T_{LT} are comparable to what is found in TmB_4 , where at $B = 0 \text{ T}$ there are two

TABLE II. Pressure dependence of ordering temperatures and Grüneisen parameter for the respective phases for uniaxial pressure along the [001] and the [110] direction as well as for hydrostatic pressure. For the continuous transition, G and E label the results obtained by exploiting the Grüneisen and Ehrenfest analysis as explained in the text. Error bars for G do not include systematic uncertainties in the overall background correction. The asterisks mark averaged values for the regimes where γ varies. Parameters for the cAFM phase are obtained above 10 K.

	dT/dp_c (K/GPa)	γ_c (mol/MJ)	$dT/dp_{\perp c}$ (K/GPa)	$\gamma_{\perp c}$ (mol/MJ)	dT/dp (K/GPa)	γ_v (mol/MJ)
T_N (G)	0.36(2)	0.65(4)	0.03(1)	0.05(2)	0.42(2)	0.76(3)
T_N (E)	0.36(9)		0.11(3)		0.6(2)	
T_{IT} (G)	-0.15*	-0.7*	0.107(13)	0.49(5)	0.04*	0.18*
T_{IT} (E)	-0.4(3)		0.10(6)			
T_{LT}	0.7(3)	1.83(13)	0.25(12)	1.01(2)	1.1(5)	3.84(10)

successive magnetic transitions, both showing positive hydrostatic pressure dependencies $dT_{N1/2}/dp \approx 0.18$ K/GPa [26]. While hydrostatic effects may often appear small due to competing uniaxial pressure dependencies canceling each other, our data also imply small uniaxial pressure dependencies, i.e., uniaxial pressure applied neither in-plane nor out-of-plane yields strong changes of the ordering temperatures. For a highly frustrated system, even small distortions are supposed to lift frustration so that large pressure effects are expected. Our data hence imply that neither of the observed long-range ordering phenomena strongly depends on magnetic frustration in this system. This also applies for the two-dimensional nature of the material since pressure along the c axis does not have a distinctly different effect on the phase boundaries than $p_{\perp c}$. The presence of multiple phase transitions is thus due to competing degrees of freedom rather than being associated with strong frustration or reduced dimensionality of the magnetic system.

The thermal expansion data do not show clear signatures associated with the CF levels which have been detected at 2.9, 12.5, and 17.7 meV [27]. Excitation of these CF levels would imply Schottky contributions to the thermal expansion in the paramagnetic phase, with a broad multiple Schottky-peak centered around 70 K which originates from the higher energy CF levels. A potential Schottky peak at around 14 K associated with the lowest CF level might be masked and shifted due to the presence of long-range magnetic order. The absence of a clear anomaly above T_N is, however, evident from Fig. 1 and the inset of Fig. 3(a). While we cannot completely rule out that the smooth changes in the data below 60 K may be partly due to CF effects, the absence of anomalies and in particular of structural anisotropy at and above the potential Schottky peak maxima at 60 and 82 K renders strong effects associated with CF excitations very unlikely. As the magnitude and the shape of the Schottky anomalies in the thermal expansion coefficients are completely determined by the energy gaps Δ_i and their pressure dependencies $\partial \ln \Delta / \partial p_i$ (see Eq. (2)), we conclude that there are rather small pressure dependencies of all Δ_i for hydrostatic pressure and uniaxial pressure along [110] and [001] (cf., [28,29]).

In order to highlight the difference of in-plane and out-of-plane behavior, a structural anisotropy parameter $\delta = (dL_{[110]} - dL_{[001]}) / (dL_{[110]} + dL_{[001]})$ is presented in Fig. 5. It represents the difference of in-plane and out-of-plane length changes, assuming finite values below $T^* \sim 65$ K and smoothly increasing cubically upon cooling to T_N . Kinks at T_N and T_{IT} , as well as a jump at T_{LT} , indicate the magnetic phase transition. While a rather conventional increase towards saturation is observed in cAFM, the distortion actually decreases in IT but eventually linearly increases in LT down to lowest accessible temperatures. The jump in length at T_{LT} implies that there is a combined structural-magnetic phase transition associated with abrupt shrinking of both $L_{[001]}$ and $L_{[110]}$, a jump-like increase of the structural anisotropy parameter δ , and evolution of the LT type of spin configuration [17,27].

The data in Fig. 5 imply that commensurate magnetic order in HT and LT promotes structural anisotropy, while in contrast δ competes with incommensurate magnetic order in IT. This competition is also evidenced by the Grüneisen parameters which, in particular, display a sign change of $\gamma_{\parallel c}$ at T_{IT} and

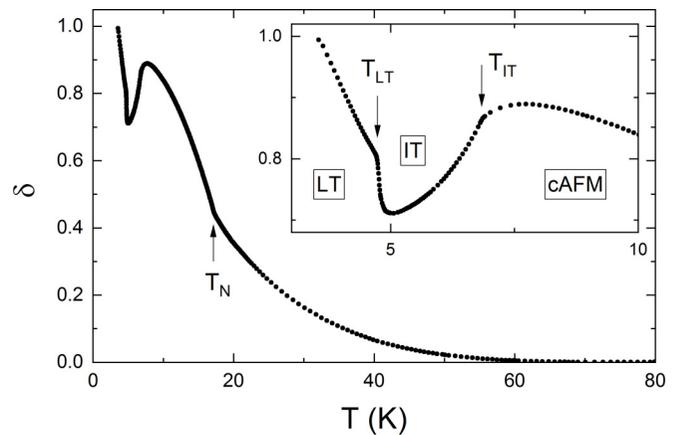


FIG. 5. Temperature dependence of the relative distortion of the uniaxial length changes $\delta = (dL_{[110]} - dL_{[001]}) / (dL_{[110]} + dL_{[001]})$. The inset shows an enlargement of the low-temperature behavior. Arrows indicate the transition temperatures of competing magnetic phases.

T_{LT} . We conclude the presence and interplay of several degrees of freedom. It is straightforward to attribute this observation to the magnetic and orbital degrees of freedom of the system. This leads to the scenario that the LT phase features combined magnetic and orbital order. This evolution is associated with the qualitative change from a volume-conserving behavior in the IT phase to a non-volume-conserving one below T_{LT} . In particular, our results do not confirm the evolution of quadrupolar order at T_N as suggested recently [18], but rather indicate a competing order parameter to evolve at T_{IT} . Note that only the magnetically incommensurate IT phase and the associated transition exhibit opposite uniaxial pressure dependencies in-plane and out-of-plane. In addition, the IT phase also shows a small but distinct non-volume-conserving temperature dependence of the Grüneisen parameters (see Fig. 4). This again indicates the presence of several degrees of freedom and may be interpreted as signature of a slowly evolving long- or short-range orbital order. Accordingly, the fluctuation regime below 10 K may indicate the presence of short-range orbital order competing with cAFM type of magnetism.

Finally, we note finite values of δ up to about 65 K, i.e., the evolution of anisotropy in the length changes as also seen in Fig. 1, indicating short-range fluctuations in this temperature regime. As visible in the inset of Fig. 1, changes in δ above T_N are associated with anomalous in-plane expansion which yields a negative thermal expansion coefficient $\alpha_{[110]}$ (cf., Fig. 2). As both magnetic and quadrupolar degrees of freedom are strongly coupled to the lattice, this phenomenon may be in general related to fluctuations of either ordering phenomena. However, it is not directly associated with fluctuations of any of the long-range ordering phenomena actually observed at low temperatures, so its nature remains unclear. Specifically, the uniaxial Grüneisen parameter $\gamma_{\perp c}$ displays a sign change at T_N which implies opposite pressure dependencies $\partial \ln \epsilon / \partial p_c$ above and below T_N , thereby excluding the short-range fluctuations being of cAFM nature. An analogous argument holds for the LT type order. We also exclude that

precursing anisotropy above T_N is associated with order of the IT type, as δ is suppressed and $L_{[110]}$ considerably shrinks in this phase.

V. SUMMARY

In summary, high-resolution capacitance dilatometry studies on the uniaxial length changes in a single crystal elucidate the interplay of spin and quadrupolar order in NdB_4 . The data imply significant magnetoelastic coupling and a structural phase transition at T_{LT} . While cAFM and LT favor structural anisotropy between in-plane and out-of-plane length changes, $\delta = (dL_{[110]} - dL_{[001]}) / (dL_{[110]} + dL_{[001]})$ competes with order in the IT phase. Notably, finite anisotropy well above T_N indicates short-range correlations which are however of neither cAFM, IT, nor LT type. Grüneisen analysis is used to derive the uniaxial as well as the hydrostatic pressure dependencies. While the ratio α/c_p features a single dominant

energy scale in LT, the data imply precursory fluctuations of a competing phase in IT and cAFM. Our results suggest the presence of orbital degrees of freedom competing with cAFM and evolving towards the magnetically and orbitally ordered ground state.

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- [1] A. P. Ramirez, *Annu. Rev. Mater. Sci.* **24**, 453 (1994).
 [2] Y. Shimizu, K. Miyagawa, K. Kanoda, M. Maesato, and G. Saito, *Phys. Rev. Lett.* **91**, 107001 (2003).
 [3] Y. Shen, Y.-D. Li, H. Wo, Y. Li, S. Shen, B. Pan, Q. Wang, H. Walker, P. Steffens, M. Boehm *et al.*, *Nature* **540**, 559 (2016).
 [4] F. L. Pratt, P. J. Baker, S. J. Blundell, T. Lancaster, S. Ohira-Kawamura, C. Baines, Y. Shimizu, K. Kanoda, I. Watanabe, and G. Saito, *Nature* **471**, 612 (2011).
 [5] S. S. Lee and P. A. Lee, *Phys. Rev. Lett.* **95**, 036403 (2005).
 [6] S. Yan, D. A. Huse, and S. R. White, *Science* **332**, 1173 (2011).
 [7] T.-H. Han, J. S. Helton, S. Chu, D. G. Nocera, J. A. Rodriguez-Rivera, C. Broholm, and Y. S. Lee, *Nature* **492**, 406 (2012).
 [8] M. J. P. Gingras and P. A. McClarty, *Rep. Prog. Phys.* **77**, 056501 (2014).
 [9] J. T. Chalker and J. F. G. Eastmond, *Phys. Rev. B* **46**, 14201 (1992).
 [10] J. S. Gardner, M. J. P. Gingras, and J. E. Greedan, *Rev. Mod. Phys.* **82**, 53 (2010).
 [11] B. S. Shastry and B. Sutherland, *Physica B* **108**, 1069 (1981).
 [12] S. Gabani, K. Flachbart, K. Siemensmeyer, and T. Mori, *J. Alloys Compd.* **821**, 153201 (2020).
 [13] R. W. Johnson and A. H. Daane, *J. Chem. Phys.* **38**, 425 (1963).
 [14] J. Etourneau and P. Hagenmuller, *Philos. Mag. B* **52**, 589 (1985).
 [15] D. Brunt, G. Balakrishnan, D. A. Mayoh, M. R. Lees, D. Gorbunov, N. Qureshi, and O. A. Petrenko, *Sci. Rep.* **8**, 232 (2018).
 [16] R. Watanuki, T. Kobayashi, R. Noguchi, and K. Suzuki, *J. Phys.: Conf. Ser.* **150**, 042229 (2009).
 [17] N. Metoki, H. Yamauchi, M. Matsuda, J. A. Fernandez-Baca, R. Watanuki, and M. Hagihala, *Phys. Rev. B* **97**, 174416 (2018).
 [18] H. Yamauchi, N. Metoki, R. Watanuki, K. Suzuki, H. Fukazawa, S. Chi, and J. A. Fernandez-Baca, *J. Phys. Soc. Jpn.* **86**, 044705 (2017).
 [19] D. Brunt, M. C. Hatnean, O. A. Petrenko, M. R. Lees, and G. Balakrishnan, *Crystals* **9**, 211 (2019).
 [20] R. Kuchler, A. Wörl, P. Gegenwart, M. Berben, B. Bryant, and S. Wiedmann, *Rev. Sci. Instrum.* **88**, 083903 (2017).
 [21] V. V. Novikov, N. V. Mitroshenkov, and A. V. Matovnikov, *J. Alloys Compd.* **646**, 906 (2015).
 [22] M. Bouvier, P. Lethuillier, and D. Schmitt, *Phys. Rev. B* **43**, 13137 (1991).
 [23] M. E. Fisher, *Philos. Mag.* **7**, 1731 (1962).
 [24] P. Gegenwart, *Rep. Prog. Phys.* **79**, 114502 (2016).
 [25] R. Klingeler, J. Geck, S. Arumugam, N. Tristan, P. Reutler, B. Büchner, L. Pinsard-Gaudart, and A. Revcolevschi, *Phys. Rev. B* **73**, 214432 (2006).
 [26] S. Gabani, G. Pristas, I. Takacova, K. Flachbart, V. Filipov, N. Shitsevalova, and K. Siemensmeyer, *Acta Phys. Pol. A* **126**, 356 (2014).
 [27] D. Brunt, PhD thesis, University of Warwick (2017), <http://wrap.warwick.ac.uk/98779/>.
 [28] K. Berggold, J. Baier, D. Meier, J. A. Mydosh, T. Lorenz, J. Hemberger, A. Balbashov, N. Aliouane, and D. N. Argyriou, *Phys. Rev. B* **76**, 094418 (2007).
 [29] J. Werner, S. Sauerland, C. Koo, C. Neef, A. Pollithy, Y. Skourski, and R. Klingeler, *Phys. Rev. B* **99**, 214432 (2019).