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Spin dynamics in bulk MnNiGa and Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn investigated by muon spin relaxationM. N. Wilson¹, T. J. Hicken¹, M. Gomilšek^{2,1}, A. Štefančič³, G. Balakrishnan³, J. C. Loudon⁴,
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We report the results of muon-spin relaxation and magnetometry investigations of bulk Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn and MnNiGa, two materials that have been proposed to host topological magnetic states in thin lamellae (antiskyrmions for Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn and biskyrmions for MnNiGa), and show spin reorientation transitions in the bulk. Our measurements reveal dynamic fluctuations surrounding the magnetic phase transitions in each material. In particular, we demonstrate that the behavior approaching the higher-temperature transitions reflects a decrease in the frequency of dynamics with increasing temperature. At low temperatures the two systems both show spin dynamics over a broad range of frequencies that persist below the respective spin reorientation transitions.

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I. INTRODUCTION

Topological magnetic objects such as skyrmions are of fundamental interest in condensed matter physics [1,2]. These states arise from a complicated hierarchy of interactions and have unique topology that brings about significant energetic stability and interesting physical effects. Recently, the known variety of such topological excitations has been expanded by the observation of two new states in thinned lamellae: biskyrmions [3] and antiskyrmions [4]. Biskyrmions are reported to consist of a bound state of two Bloch skyrmions with the same chirality, and hence have a topological charge of $N = 2$, while antiskyrmions are objects where the winding varies from Néel type to Bloch type around the circumference of the object, and have a topological charge of $N = -1$ [2].

Biskyrmions have been reported in the layered manganite La_{2-2x}Sr_{1+2x}Mn₂O₇ [3] and in certain compositions of hexagonal MnNiGa [5,6], while antiskyrmions have so far been reported in Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn [4] and related centrosymmetric Heusler systems [7,8]. Indirect evidence for the presence of a topologically nontrivial state has been reported for these materials, and other Heusler materials related to the antiskyrmion hosts [9–12], however direct evidence for both the biskyrmion and antiskyrmion states has thus far relied on Lorentz transmission electron microscopy (LTEM). However, this identification is controversial, as further LTEM and x-ray holography measurements of MnNiGa suggest that the LTEM images can be explained as reflecting a more conventional type-II magnetic bubble state [13], while theoretical studies continued to support the existence of biskyrmions in centrosymmetric magnetic films [14]. Nevertheless, these materials remain of interest as a result of their complex magnetic behavior.

Both MnNiGa and Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn undergo two magnetic transitions with decreasing temperature [4,6], with a ferromagnetic state occurring below $T_C = 400$ (Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn) or 350 K (MnNiGa), and a lower-temperature transition below 250 K in both materials that is accompanied by a change in magnetic moment. In Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn, the transition is suggested to be a spin reorientation by analogy with similar materials such as Mn_{1.4}PtSn [4,15] and Mn₂RhSn [16]. Neutron diffraction measurements on MnNiGa show that the low-temperature transition in this material also involves spin reorientation, in this case introducing an antiferromagnetic component to create a canted noncollinear state [17,18].

Topological Hall effect anomalies have been reported in both materials. Below T_C in 50- μ m-thick polycrystalline samples of MnNiGa, from $0.2 < \mu_0 H < 1$ T, a topological Hall signal is observed, consistent with the region in which the biskyrmion state is reported in sub- μ m-thick lamellae [6]. In contrast, in Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn a topological Hall signal is only reported below the spin-reorientation transition for $T < 150$ K [8] in bulk samples. This suggests that some sort of topologically nontrivial state exists in bulk samples of both materials, and not only in thin films or lamellae.

In this paper we present a study of the magnetic behavior of bulk MnNiGa and Mn_{1.4}Pt_{0.9}Pd_{0.1}Sn using a combination of magnetometry and muon-spin relaxation (μ SR). Our measurements reveal dynamic fluctuations around the two transitions in these materials that appear as a function of temperature. Furthermore, the μ SR underscores the difference in the nature of the fluctuating magnetism occurring between the lower temperature spin reorientation transitions. The investigation of these fluctuations sheds light on the bulk behavior of

these materials and helps put the observation of topologically nontrivial magnetic states in thin lamellae in context.

II. METHODS

We performed muon-spin relaxation (μ SR) and magnetometry measurements on samples of $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ and MnNiGa [nominally $(\text{Mn}_{0.5}\text{Ni}_{0.5})_{65}\text{Ga}_{35}$] synthesized at the University of Warwick, UK. These samples consist of polycrystalline boules prepared by arc-melting the components together. The $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ boule is known to have regions of twinned crystallites that might introduce a preferred orientation to the sample, rather than having a fully random distribution of orientations.

For the $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ μ SR measurements, we used two halves of the boule with circular cross section, weighing a total of 1.217 g, with a total cross-sectional area of approximately 0.9 cm^2 , while for the MnNiGa measurements we used a single boule of mass 0.907 g and circular cross-sectional area of 0.6 cm^2 . μ SR measurements were performed on the HiFi spectrometer at the STFC-ISIS pulsed muon source. The MnNiGa measurements were performed in two separate experimental runs, with the 0.01 and 0.3 T measurements conducted first, and the other measurements later. Small differences in the setup of these experiments led to a difference in the observed high-temperature relaxing asymmetries. For the magnetometry measurements we cut an 8.7 mg piece from the $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ boule and a 4.8 mg piece from the MnNiGa boule, and measured them using a Quantum Design MPMS. AC magnetometry measurements were collected with a drive field of 0.3 mT at 111 Hz for $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ and 0.35 mT at 113 Hz for MnNiGa .

III. RESULTS AND DISCUSSION

A. MnNiGa

1. Magnetometry

Figure 1 shows magnetic field-temperature phase diagrams for MnNiGa collected from (a) zero-field cooled (ZFC) temperature scans of DC magnetization (plotted as magnetization divided by field M/H), and (b) the real part of the AC susceptibility (χ'). (Examples of these measurements are shown in Fig. 2.) The DC magnetization shows a transition at $T_C = 340\text{ K}$, characterized by a steplike increase in M [also evident in, e.g., Fig. 2(b2)], which remains nearly unchanged with applied magnetic field H . The AC magnetometry also shows a feature at this temperature that does not significantly change with H , showing up as a kink or weak maximum in χ' (most obviously at higher field values). This sort of feature has been seen in other materials such as the skyrmion host Cu_2OSeO_3 , and interpreted as the transition from a field polarized state to a high temperature paramagnetic state [19]. There is also a second high-temperature feature in χ' that is not resolved in the M/H data. This shows up as a sharp steplike increase in χ' that rapidly decreases in temperature as the field is increased, reminiscent of the transition from the conical to field polarized states in chiral helimagnets [19]. Below 200 K, the data show a small change in both magnetization and χ' , suggesting a second magnetic transition occurring at all magnetic fields,

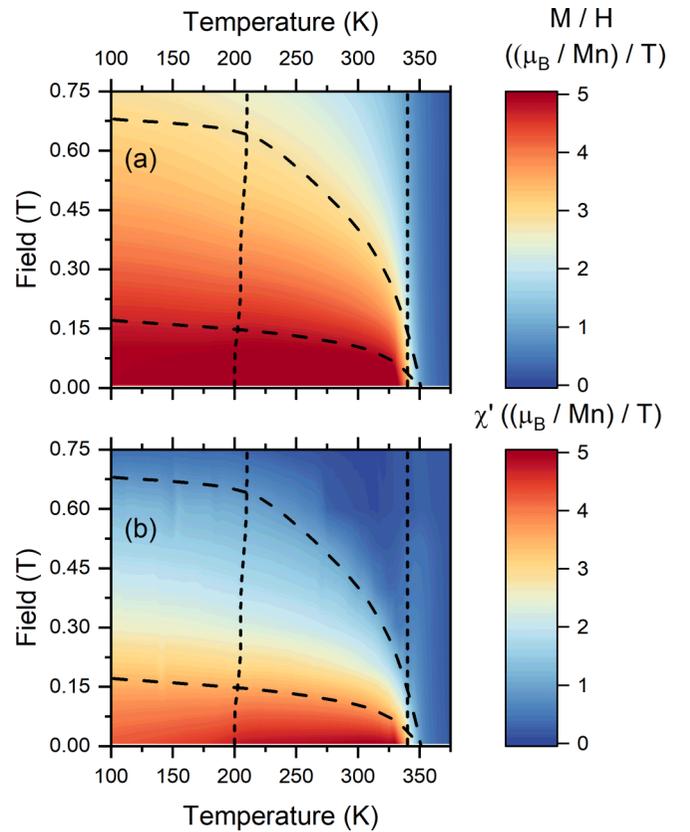


FIG. 1. Phase diagrams for MnNiGa measured through fixed magnetic field temperature sweeps on warming, after cooling in zero field. (a) DC magnetization, plotted as magnetization divided by applied field; (b) real part of the AC susceptibility (χ'). The lines show transitions determined by DC magnetometry temperature scans (short dashed lines, maxima in dM/dT for low-temperature transition, minima for the higher temperature transition), and AC magnetometry (long dashed lines, minima in $d\chi'/dH$).

consistent with previous reports. The transition temperature is determined from a peak in dM/dT (see Fig. 3).

Figure 2(a) shows the magnetic field derivative of the real part of the AC susceptibility $d\chi'/dH$ at 230 K which has two pronounced minima at around 0.12 and 0.62 T (data at other temperatures show similar behavior). This suggests that there is a crossover in magnetic behavior as a function of field, rather than a smooth evolution. The minimum at 0.62 T likely corresponds to the transition to a field polarized state, while the origin of the low-field transition is less clear. A similar series of two transitions was seen in MnNiGa in topological Hall effect measurements [6], and reported to arise from a low-field stripe/helical state, with biskyrmions reported at intermediate fields between the two transitions. Although, these states have only been reported in thin lamellae [6], our observation of a similar phase diagram in bulk samples suggests that these transitions may persist in the bulk.

2. μ SR

To investigate the local magnetism in MnNiGa , we measured longitudinal field μ SR. In this technique, spin polarized muons (with spins initially antiparallel to the longitudinal field

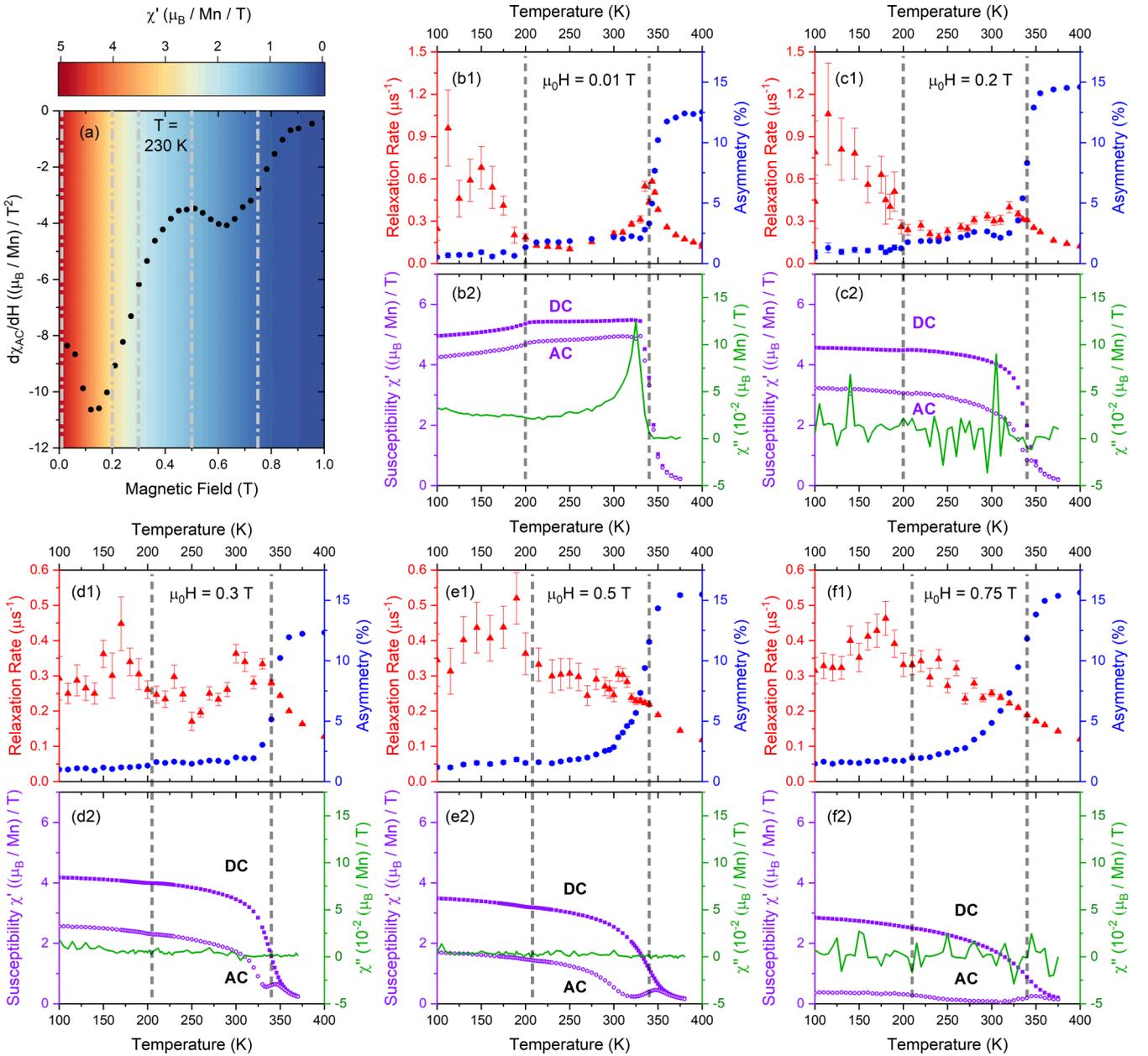


FIG. 2. Magnetometry and μ SR data for MnNiGa measured at several magnetic fields as a function of temperature after zero-field cooling. (a) Derivative of the real part of the AC susceptibility $d\chi'/dH$ (black circles), as a function of field after ZFC to 230 K, and χ' as the color map. (b)–(f) subpanel (1): Measured μ SR relaxation rate (λ , red triangles, left axis), the relaxing asymmetry (A_r , blue circles, right axis); and subpanel (2): DC magnetization plotted as M/H (in purple squares, left axis), χ' (purple open circles, left axis), and the imaginary part of the AC susceptibility (χ'' , green solid line, right axis). Gray dashed lines in (b)–(f): Temperatures of magnetic transitions determined from peaks in dM/dT at each magnetic field (maxima for the low-temperature transition, minima for the higher temperature transition). The white dashed lines in (a) indicate the magnetic fields where panels (b)–(f) were measured.

B_L) are implanted into the material, where they then precess in the local magnetic field before decaying. The average spin polarization of the muon ensemble is determined by measuring the asymmetry function $A(t) = [F(t) - B(t)]/[F(t) + B(t)]$, where F represents the number of decay positrons measured forward of the initial muon spin direction and B is the number measured in the backwards direction. In an ordered magnet there are typically two principal contributions to the time evolution of the asymmetry: (i) a contribution arising from

muons coherently precessing in the internal field directed perpendicular to the initial muon spin direction, and (ii) a contribution arising from muon spins initially parallel to the local magnetic field dephased by time-dependent fluctuations. The first of these contributions would often result in damped oscillations in the asymmetry. However, the time resolution of a pulsed muon source limits our ability to resolve these features, and this contribution gives rise to a fraction of “missing asymmetry” that relaxes too quickly to resolve. In contrast,

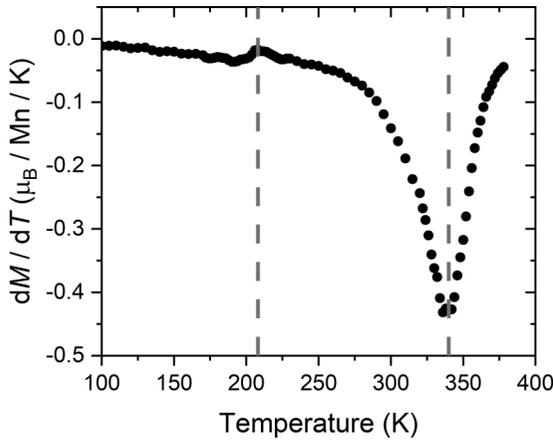


FIG. 3. dM/dT for MnNiGa measured at 0.5 T. Gray dashed lines indicate transitions at 208 and 340 K determined from the extrema.

the second contribution results in a slow decay of the asymmetry, determined by the spectral density $S(\omega)$ of local field fluctuations in the material, which is a Fourier transform of the local field-field autocorrelation function $\gamma_\mu^2 \langle \delta B(\tau) \delta B(0) \rangle$, where $\gamma_\mu = 2\pi \times 135.5$ MHz/T is the muon gyromagnetic ratio and δB is the fluctuating field at the muon site. The fluctuations will relax the muon spins most effectively when there is spectral density at the frequency associated with the applied longitudinal field $\omega_0 = \gamma_\mu B_L$ (≈ 8 – 600 Mrad/s in this work, depending on the applied field B_L), giving a longitudinal relaxation rate in the long-time limit of

$$\lambda = \gamma_\mu^2 \int_0^\infty d\tau \cos \omega_0 \tau [\langle \delta B_x(\tau) \delta B_x(0) \rangle + \langle \delta B_y(\tau) \delta B_y(0) \rangle], \quad (1)$$

where δB_x and δB_y represent field fluctuations perpendicular to the direction of the applied field B_L [20].

Figure 4 shows representative μ SR data for the MnNiGa sample, measured in a longitudinal field of 0.01 T. The behavior is qualitatively similar at other measured fields.

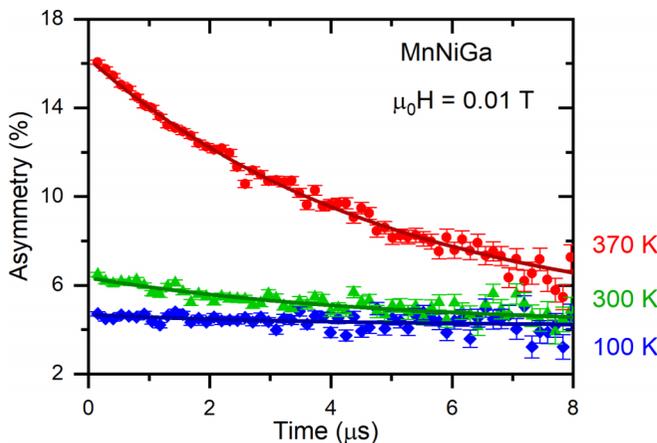


FIG. 4. Representative μ SR data for MnNiGa measured in a longitudinal field of 0.01 T at 100 K (blue), 300 K (green), and 370 K (red). Solid points show measured data; lines show fits to Eq. (2).

At each temperature the data can be well described by a single exponentially relaxing component, which relaxes to a temperature-independent baseline (4.2% for the field shown). This behavior is characteristic of muons dephased by dynamic fluctuation of magnetic moments in the material, and therefore provides information about how the dynamics vary across the phase diagram. We fit these data to

$$A(t) = A_b + A_r \exp(-\lambda t), \quad (2)$$

where A_b is the baseline asymmetry, A_r is the relaxing asymmetry, and λ is the relaxation rate. The results of this fitting are shown in Fig. 2, along with the corresponding temperature scans from magnetometry at equivalent magnetic fields. Dashed lines in Fig. 2(a) show the magnetic fields where we measured μ SR spectra, showing that we have one measurement in the low-field state, three spanning the intermediate field region, and one in the field-polarized state.

The data measured at 0.01 T in Fig. 2(b) show a peak in the μ SR relaxation rate λ at 340 K, along with a steplike drop in the relaxing asymmetry typical of a transition to long-range order. These features occur at the same temperature as the minimum in dM/dT that marks T_C in the magnetometry data. In the fast-fluctuation limit, the μ SR relaxation rate λ is often described by the Redfield formula [20],

$$\lambda = \frac{2\Delta^2\nu}{\omega_0^2 + \nu^2}, \quad (3)$$

where $\Delta = \gamma_\mu^2 \sqrt{\langle \delta B^2 \rangle}$ is determined by the amplitude of the fluctuations and ν is a fluctuation rate. If the relaxation is determined by fluctuations with a characteristic fluctuation rate ν that decreases as the transition temperature is approached from below, then the μ SR relaxation rate λ would be expected to show a peak slightly below T_C which broadens and shifts to lower temperature as the applied magnetic field is increased, as illustrated by Fig. 5 [21]. Our data follows this general trend of behavior, with the μ SR data at larger magnetic fields showing a high temperature peak (near T_C) in the relaxation rate that is broader, weaker, and shifted to lower temperatures compared to the 0.01 T data.

However, a potentially puzzling feature occurs in the imaginary part of the AC magnetic susceptibility [χ'' , green line, lower panel of Fig. 2(b)], which shows a sharp peak 15 K lower in temperature than the peak in the μ SR relaxation for the 0.01 T data set. AC susceptibility probes a substantially lower frequency range than μ SR, and therefore we might expect the influence of fluctuations with a single characteristic fluctuation rate that decreases with increasing T to appear at higher temperatures in AC than in the μ SR. The observation of a peak in χ'' at lower temperature therefore suggests that there are further spin fluctuations leading to additional dynamics at frequencies that are too low to be seen with μ SR.

Furthermore, there is no obvious signature of the high temperature transition in χ'' at any of the fields above 0.01 T. It is possible that this simply arises from a broadening of the transition with field, which would reduce the measured peak height of the imaginary signal, potentially bringing it below the noise floor. However, since Fig. 2(a) indicates that at these higher fields the system is in a different magnetic state, it may instead be that the low-frequency spectral weight that gives

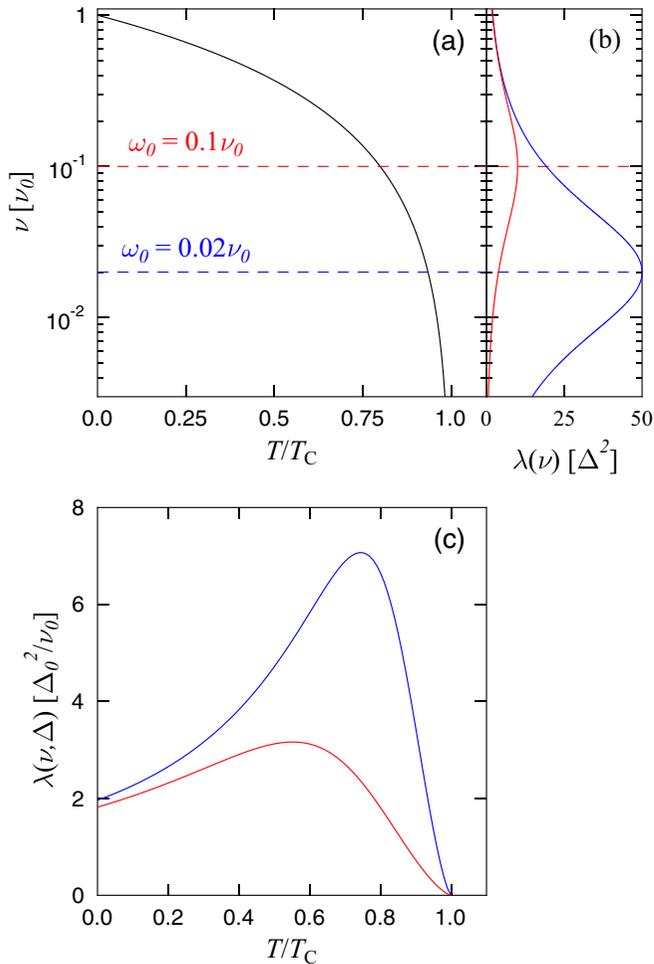


FIG. 5. Schematic showing the expected behavior of the μ SR relaxation rate λ from a fluctuation with a frequency ν that decreases approaching T_C from below. (a) Fluctuation rate as a function of temperature, modeled using a 3D Heisenberg power law $\nu = \nu_0(1 - \frac{T}{T_C})^{\zeta z}$, where ν_0 is the zero-temperature frequency, $\zeta = 0.7048$ is the correlation-length critical exponent, and $z = 2.035$ is the dynamical critical exponent [22]. (b) Relaxation rate (x axis) as a function of frequency (y axis) following the Redfield formula, for the case of constant Δ , for two different values of the applied field ($\omega_0 = \gamma_\mu B_L$). (c) Relaxation rate as a function of temperature, also now including a power-law decrease in Δ with increasing temperature [21]. Red line is for $\omega_0 = 0.1\nu_0$, blue line is for $0.02\nu_0$.

rise to the peak in χ'' at 0.01 T does not persist into these higher magnetic field states.

At lower temperature in the μ SR data, we see a further drop in the relaxing asymmetry at 200 K, which coincides with a maximum in dM/dT . This is consistent with the temperature of the spin-reorientation transition to a noncollinear canted ferromagnetic state that has been suggested from neutron scattering measurements [18]. The additional drop in the relaxing asymmetry would, for example, occur if a larger fraction of the static magnetic moments in the material are canted away from the applied field direction below 200 K, which would increase the fast-relaxing component of the μ SR signal, leading to the additional missing asymmetry we observe. This transition does not coincide with a peak in the

μ SR relaxation rate, nor a peak in the imaginary part of the magnetic susceptibility. Instead, there is a broad increase in both of these quantities below the 200 K transition, suggesting that the low-temperature dynamics occur over a wide range of frequencies (from 10s of Hz to MHz) and persist across a broad temperature range. As alluded to by Bogdanov *et al.* [23], this situation can arise from a near degeneracy in canting angles around the field direction, resulting in a small energy barrier between magnetic domains that would therefore undergo significant dynamics.

We see limited qualitative difference in the μ SR behavior crossing from 0.2 to 0.3 T, where the previous reports suggest a transition from the helical or stripe state to the biskyrmion state. We also do not see evidence for a transition between these states as a function of temperature, as might be expected to occur between 250 and 350 K for the 0.2 T data set. This suggests that the proposed transition from the helical to biskyrmion states does not produce a sizable enough change to the local magnetic fields or their dynamics to give a resolvable response in the μ SR data, which contrasts with previous μ SR measurements of Bloch and Néel type skyrmions [21,24,25], where the presence of skyrmions is detectable. Therefore, our μ SR data suggest that either the biskyrmion state does not host characteristic dynamics similar to those of conventional skyrmions, or this biskyrmion state does not persist in bulk samples, as might be expected if the interpretation of biskyrmions as type-II bubbles stabilized by demagnetization and confinement effects is correct [13,26].

B. $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$

1. Magnetometry

We now turn to discussion of the data measured on the reported antiskyrmion host $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$. The DC magnetization in Fig. 6 shows two features at each field: a sharp rise with decreasing temperature occurring around 390 K, and a second rise starting at 145 K. This indicates the presence of two separate magnetic transitions in the material, at $T_C \approx 380$ K and $T_{SR} \approx 140$ K, consistent with previous reports [4]. The rise in magnetization at the lower-temperature transition, in contrast to MnNiGa, shows that the spin-reorientation transition involves an increase in the ferromagnetic component of the magnetism at lower temperature. The AC susceptibility also shows two features coinciding with the two transitions, although these are slightly different at the different magnetic fields. At 0.1 and 0.3 T, T_C is identified by a small peak in χ' with no feature in χ'' , while T_{SR} shows up as a steplike drop in χ' and a peak in χ'' . At 0.01 T, both transitions show a peak in χ'' . This indicates that the lower temperature transition shows dynamics with frequencies close to the AC drive field frequency of 111 Hz at all fields. At temperature below T_{SR} , χ'' continues to increase down to low temperature in the 0.1 and 0.01 T data, which suggests that there are spin dynamics at low T that are at least partially suppressed by the application of a 0.3 T magnetic field.

2. μ SR

Figure 7 shows representative μ SR data measured in a longitudinal magnetic field of 0.1 T at several temperatures.

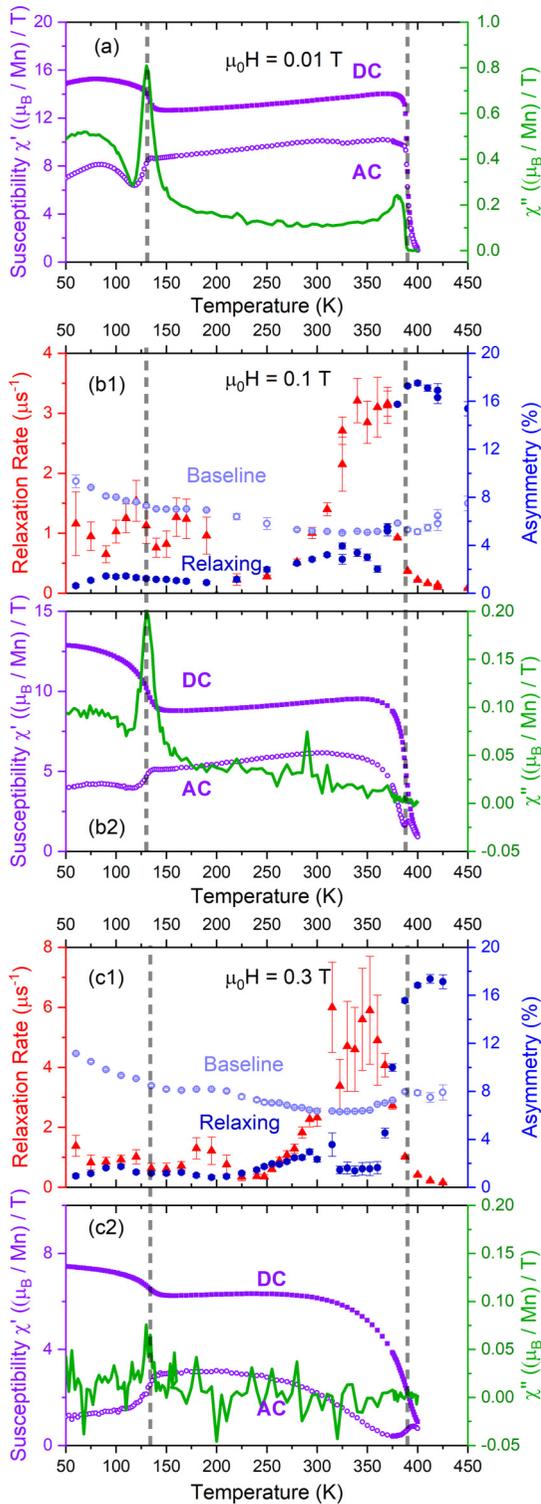


FIG. 6. (a) Magnetometry data for $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ measured at 0.01 T. μSR and magnetometry data measured at (b) 0.1 T and (c) 0.3 T after zero-field cooling. Panel (1) shows μSR relaxation rate (red triangles, left axis), baseline asymmetry (open light blue circles, right axis), relaxing asymmetry (blue circles, right axis). Panel (2) shows DC magnetization plotted as M/H (purple squares, left axis), real part of the AC susceptibility (χ') (purple open circles, left axis), and imaginary part of the AC susceptibility (χ'') (green solid line, right axis). Gray dashed lines indicate transition temperatures determined from minima in dM/dT at each magnetic field.

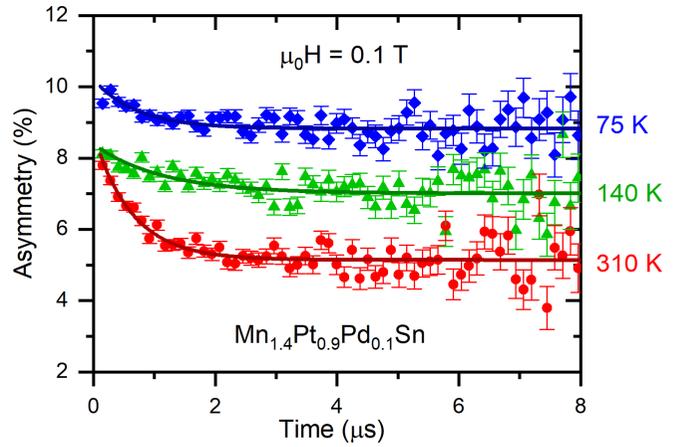


FIG. 7. Representative μSR data measured in longitudinal fields of 0.1 T for $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ at 75, 140, and 310 K. Solid points show measured data; lines show fits to Eq. (2).

These data again show exponential relaxation that can be modeled by Eq. (2), with the notable difference compared to the MnNiGa data that the nonrelaxing baseline (A_b) in these data varies as a function of temperature. Baseline asymmetry varying with temperature is unusual in μSR , as this term is normally temperature independent and arises from muons hitting the silver sample holder (which show no relaxation in a longitudinal field greater than about 5 mT). We therefore fit the μSR data to Eq. (2), allowing the baseline, relaxation rate, and relaxing asymmetry to freely vary with temperature. The results of this fitting procedure are shown in Fig. 6 for fields of 0.1 and 0.3 T, along with magnetometry data at the same magnetic fields for comparison, and magnetometry data measured at 0.01 T as a reference.

We see evidence for the two magnetic transitions in the μSR data shown in Figs. 6(b1) and 6(c1). The higher-temperature transition is characterized by a large peak in the relaxation rate at around 360 K, along with a sharp drop in the relaxing asymmetry. This peak is centered ≈ 30 K below T_C , and is not accompanied by a peak in χ'' . This broad peak suggests that there are dynamics in the system below T_C with a fluctuation rate that decreases slowly in frequency as the temperature is increased towards T_C , hence putting its characteristic frequency within the μSR time window over a broad temperature range, similar to the picture shown in Fig. 5.

It has been reported from magnetic entropy measurements that bulk $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ hosts antiskyrmions up to 350 K or higher at 0.1 T, while it does not host them until below 300 K at 0.3 T [27]. It is plausible that near the higher-temperature transition we are probing dynamics arising from characteristic excitations of the antiskyrmion state at this lower field. However, this would require that the antiskyrmion state is characterized by a lower relaxation rate than the neighboring magnetic states, in contrast with published data on Bloch and Néel skyrmions, where an increase in this quantity is seen [21,22,24,25].

At lower temperatures, there are further features in the μSR data: two small peaks in the relaxation rate, around 175 and 100 K for both fields, which roughly correspond with

kinks in the baseline asymmetry, and a rise in the relaxation rate at the lowest measured temperatures which is correlated with a further drop in the relaxing asymmetry. The behavior of the relaxing asymmetry in particular is quite distinct to that in MnNiGa, showing that these transitions behave differently, despite both being spin reorientations, with the complicated temperature dependence in $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ around the low-temperature transition suggesting a multistep evolution of the dynamics of the magnetism at these temperatures.

Furthermore, the continued rise in the imaginary part of the susceptibility below T_{SR} , and enhancement in the μSR relaxation rate at the lowest temperatures, suggests that there are low-temperature spin dynamics across a range of frequencies in this low temperature phase, similar to what was seen in MnNiGa. Recent reports have suggested that there is a complicated evolution between different geometries of bubbles or antiskyrmions as a function of temperature and magnetic field in thin plates of $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ [28,29], with circular bubbles, elliptical bubbles, triangular bubbles, and square “antiskyrmions” all appearing at different points of the phase diagram. Transitions between these different states as a function of temperature may explain the multistep evolution of the dynamics indicated by our μSR data.

The other notable feature at low temperatures is the continued rise in the baseline asymmetry with decreasing temperature, with a plateau approximately between the two low temperature peaks in the relaxation rate. This feature likely arises from a combination of two factors. First, if we have some preferred orientation in our polycrystalline sample (as is likely), and the spin reorientation transition causes the internal magnetic field to more frequently point along the initial muon polarization direction due to this preferred orientation, there will be a reduction in the static contribution to the muon signal (which is seen as missing initial asymmetry in our data). This would result in an increasing initial asymmetry, as is seen between the 140 and 75 K data sets in Fig. 7. However, in this case, regions of the sample where the internal field lies parallel to the muon spin would still show dynamic relaxation caused by fluctuations of the moments. Our data shows both a nonrelaxing contribution to the muon asymmetry, and a contribution with relaxation arising from dynamics. This suggests that there are spatially separated regions of the sample that show different dynamics: regions where the dynamics are either too fast or too slow to be probed by μSR (giving a nonrelaxing signal), and regions where dynamics fall into the μSR frequency window (giving a relaxing signal). Since the total change in the baseline asymmetry with temperature is approximately 5%, and the maximum relaxing asymmetry is 20%, we can infer that these regions with dynamics outside the μSR frequency window occupy approximately a quarter of the sample volume. It is possible that these different regions arise from effects near structural crystal domain boundaries. Notably, this material is known to be prone to twinning defects [4], and transmission electron microscopy (TEM) images of thin lamella taken from of our sample (see Fig. 8) demonstrate that we have closely spaced (<50 nm) structural twinning defects throughout our sample. Magnetism varying slightly surrounding these defects (as previously reported [4]) would result in a substantial portion of the sample experiencing different dynamics, explaining our observations.

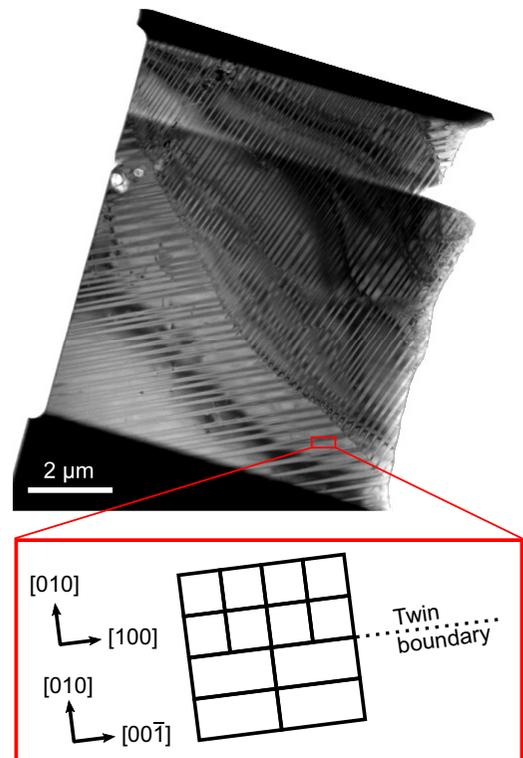


FIG. 8. Bright field TEM image of $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ showing the presence of twinning domains in the crystal structure.

IV. CONCLUSION

In conclusion, we have presented μSR and magnetometry of bulk samples of two materials (MnNiGa and $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$) thought to host skyrmionic states in thin lamellae form, which show bulk spin reorientation transitions with temperature. In both samples we find evidence for two magnetic transitions as a function of temperature. Our data suggest that the higher-temperature transition in both samples is characterized by dynamics whose fluctuation rate decreases in frequency with increasing temperature. In both materials we also find evidence for spin dynamics persisting down to low temperatures below the spin-reorientation temperatures. However, the spin reorientation transitions themselves show rather different signatures in the two systems, with $\text{Mn}_{1.4}\text{Pt}_{0.9}\text{Pd}_{0.1}\text{Sn}$ showing an increase in the ferromagnetic moment and the peak in χ'' but also suggesting some spatial variation of dynamics across the sample from μSR , while the MnNiGa data suggest a more continuous change in dynamics across this transition. Finally, we note that the μSR measurements do not show unambiguous evidence for dynamics reflecting the proposed biskyrmion state in MnNiGa, which would suggest that either this state does not occur in bulk samples, or that its dynamics are different from those seen in other skyrmionic systems, and are not distinguishable from those of the low-temperature helical or stripe phase. However, this work does further advance our understanding of the behavior of the bulk materials, which underpins the observation of topological magnetism in thin lamellae.

Data presented in this paper are available at [30].

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