Simultaneous Measurement of Proton and Lepton Kinematics in Quasielasticlike $\nu_\mu$-Hydrocarbon Interactions from 2 to 20 GeV


(Received 23 March 2022; revised 25 May 2022; accepted 21 June 2022; published 6 July 2022)

Neutrino charged-current quasielastic-like scattering, a reaction category extensively used in neutrino oscillation measurements, probes nuclear effects that govern neutrino-nucleus interactions. This Letter reports the first measurement of the triple-differential cross section for $\nu_\mu$ quasielastic-like reactions using the hydrocarbon medium of the MINERvA detector exposed to a wideband beam spanning $2 \leq E_\nu \leq 20$ GeV. The measurement maps the correlations among transverse and longitudinal muon momenta and summed proton kinetic energies, and compares them to predictions from a state-of-art simulation. Discrepancies are observed that likely reflect shortfalls with modeling of pion and...
nucleon intranuclear scattering and/or spectator nucleon ejection from struck nuclei. The separate determination of leptonic and hadronic variables can inform experimental approaches to neutrino-energy estimation.

DOI: 10.1103/PhysRevLett.129.021803

Current and future long baseline neutrino experiments [1–4] seek to delineate the neutrino mass ordering and to quantify the presence of charge-parity violation in the neutrino sector. These experiments will use neutrinos of energies from 0.3 to 4 GeV and higher if tau-neutrino appearance is explored [5]. Accurate models of neutrino-nucleus interactions are required to relate the energies of visible final-state particles of events observed in the detectors to the initiating true neutrino energies that underwrite the oscillations of neutrino flavor. A leading contributor to charged-current neutrino interactions at these energies is the quasielastic-like channel:

\[ \nu_\mu + A \rightarrow \mu^- + \text{nucleons} + A'. \]  

Charged-current neutrino interactions within nuclei, even those with apparent two-body quasielastic final states, are altered by a number of poorly understood effects: The struck nucleons of the initial state are bound and in motion [6,7]; short-range multinucleon processes give rise to enhanced reaction rates relative to scattering on free nucleons [8–13], and hadrons produced in the parent \( \nu_\mu \) interactions with nucleons undergo intranuclear final-state interactions (FSI) within the target nuclei. While the reaction \( \nu_\mu + A \rightarrow \mu^- + p + A' \), wherein nearly all final-state energy is visible, is thought to be the main contributor to the quasielastic-like channel [Eq. (1)] a significant number of events may have energy deposited in undetected neutrons or light nuclear fragments. Final states of the latter kind complicate the task of inferring neutrino energy from samples of quasielastic-like events.

The reaction of Eq. (1) has received repeated experimental scrutiny; however only single- or double-differential cross sections in muon kinematics have been reported, mostly carried out with \( \nu_\mu \) of incident energies, \( E_\nu \), of sub-GeV to few GeV [14–31]. This Letter reports a new measurement of the quasielastic-like channel in which the final-state muon transverse \( (p_t) \) and longitudinal \( (p_\parallel) \) momenta are measured in each event simultaneously with the total “available” (calorimetrically visible) recoil energy \( (E_{\text{available}}) \) used in previous analyses of data from MINERvA [13,32,33]. Since the signal requires final state muon plus nucleons only, \( E_{\text{available}} \) is the sum of the kinetic energies of all protons, denoted \( \Sigma T_p \).

Under the assumption of a stationary target neutron in \( \nu_\mu + n(\text{bound}) \rightarrow \mu^- + p \), energy transfer also can be inferred to be

\[ d_0^{(\text{QE})} \equiv \frac{m^2_p - (m_n - E_b)^2 - m^2_\mu + 2(E_\mu - p_\mu \cos \theta_\mu)E_\mu}{2(m_n - E_b) - E_\mu + p_\mu \cos \theta_\mu}. \]  

Here, \( m_\mu, m_p, \) and \( m_n \) are the masses of the muon, proton, and neutron, \( E_b \) is the average binding energy of 34 MeV [6,34,35], and \( E_\mu, p_\mu, \) and \( \theta_\mu \) are the muon energy, momentum, and angle with respect to the neutrino beam. The \( d_0^{(\text{QE})} \) of Eq. (2) is the quantity added to the reconstructed muon energy by T2K to estimate neutrino energy of quasielastic-like events, while \( \Sigma T_p \) is the amount added to the muon energy by NOvA to form its neutrino energy estimator for events with quasielastic-like topologies. Combined T2K-NOvA analyses will be credible to the extent that interaction models correctly predict the relationship between these quantities. The measurements of this Letter elicit the correlations among the kinematic variables of quasielastic channels, thereby confronting the models with information of a kind that heretofore has not been available.

The analysis uses high-statistics samples of \( \nu_\mu \) charged-current interactions recorded by the MINERvA detector [36] exposed to the wideband, medium energy NuMI beam [37] at Fermi National Accelerator Laboratory. In the NuMI beam, 120 GeV protons impinging upon a carbon target produce pions and kaons that are subsequently charge-selected, focused by a magnetic horn system, and directed into a pipe where they decay. The resulting neutrino flux is calculated using a \textsc{geant4} simulation of the beam optics with input from hadronic interaction data relevant to the beam and materials [38]. The neutrino flux is constrained by previous measurements of neutrino elastic scattering from atomic electrons, \( \nu e^- \rightarrow \nu e^- \) [39]. This constraint reduces the normalization uncertainty from 7.8% to 3.9% for muon neutrinos of energies between 2 and 20 GeV. The neutrino interactions occur in the central scintillator tracker of the MINERvA spectrometer, which has a mass fraction of 88.5% carbon, 8.2% hydrogen, 2.5% oxygen, and trace amounts of other elements. Primary vertices of selected events are restricted to a central 5.3 ton region. The spatial resolution of the tracker enables reconstruction of final-state protons and Michel electrons from the \( \pi^+ \rightarrow \mu^+(\nu_\mu) \rightarrow e^+(\nu_e \bar{\nu}_\mu) \) decay chain, as well as the tracks of muons. The magnetized MINOS near detector, located downstream of MINERvA, is used to determine the charge and momenta of exiting muons. The scintillator tracker and the surrounding sampling calorimeters enable calorimetric
Occasionally, final-state neutrons leave a small amount of energy that is tagged as a photon or included in $\Delta T_p$; the reference simulation predicts and corrects for this effect. The average $\Delta T_p$ for protons is $\approx 250$ MeV; neutrons contribute less than 10 MeV of energy in 74% of events and an average of 85 MeV for the rest.

The MINERvA detector response is simulated using GEANT4 [40] version 4.9.4p2 with the QGSP_BERT physics list. The optical and electronics performance is also simulated. Through-going muons are used to determine the absolute energy scale. Full descriptions of calibrations are given in Refs. [36,41]. The absolute energy response to charged hadrons is set according to measurements using a charged particle test beam [42] and a scaled-down version of the MINERvA detector. The effects of accidental activity as a function of beam intensity are simulated by overlaying hits from data in both MINERvA and MINOS.

The reference signal and background models for this analysis are based on a modified version of the GENIE [43] v.2.12.6 event generator. Quasielastic interactions are modeled using the Llewellyn Smith formalism [44] with an axial-vector form factor [45] and an axial-vector form factor based on a $z$-expansion fit to deuterium data [46]. Resonance production is simulated using the Rein-Sehgal model [47] with a dipole axial mass of $M_A^{\text{RES}} = 1.12$ GeV/c$^2$. The nuclear initial state is a relativistic Fermi gas model [48] with $k_F = 0.221$ GeV/c and with a Bodek-Ritchie high momentum tail [49]. Multinucleon quasielastic-like interactions are simulated by the “Valencia model” described in Refs. [10,11,50]. Intranuclear final-state interactions of produced hadrons are modeled using the INTRANUKE-HA package [51].

To better describe MINERvA data, a number of modifications are made in the reference model that are collectively denoted MINERvA tune v4.4.1. The quasielastic cross section is modified as a function of energy and three-momentum transfer based on the random phase approximation of the Valencia model [52,53] appropriate for a Fermi gas [54,55] to account for long-range correlations between nucleons. To account for an observed excess in specific regions of three-momentum transfer and $\Delta T_p$, the multinucleon cross section is increased based on fits to MINERvA data [13] from a lower energy beam configuration. Additionally, based on fits to $\nu_\mu$-hydrogen data [56], the nonresonant charged-current pion production is decreased by 43%, the overall baryon-resonance pion production is increased by 15%, and $M_A^{\text{RES}}$ is set to 0.94 GeV.

Samples for measuring quasielastic-like interactions and their backgrounds require a muon track that starts in the fiducial volume and is identified in MINOS as negatively charged. All other tracked particles originating from the interaction vertex at the beginning of the muon track must have $dE/dx$ consistent with a proton. Signal and background samples are formed by counting the number of Michel electron candidates within 600 mm long, 600 mm diameter cylinders centered on the neutrino vertex and on endpoints of tracked particles, and by counting isolated clusters constructed from two-dimensional clusters with at least 1 MeV visible energy. The former identify $\pi^+$, and the latter identify photons from $\pi^0$ decays. Clusters with an

![FIG. 1. The flux-averaged triple-differential cross section for quasielastic-like events, $\frac{d^3\sigma}{dp_T dp_T dp_T}$, shown as points with colored error bands for designated intervals of $p_T$ in panels of $p_T$. Note the use of scaling factors and log scale to elicit the trends and consistency across all $p_T$. The predictions of the reference model MINERvA tune v4.4.1 are shown as lines in each panel.](image)
energy less than 10 MeV per hit are assumed to be caused by neutrons producing low energy protons and are not used. Events that contain either a $\pi^+$ ($67\%$), a $\pi^0$ ($19\%$), or both ($14\%$), comprise the dominate backgrounds to the quasielastic-like signal. Four exclusive samples are assembled using the criteria of $0$ or $\geq 1$ Michel electrons, and $\leq 1$ or $\geq 2$ isolated clusters. Sample A with no Michel electrons and $\leq 1$ isolated cluster is the signal sample. Sample B has a Michel electron but $\leq 1$ isolated cluster and is rich in single $\pi^+$ events. Sample C comprises events with $\geq 2$ isolated clusters but no Michel electrons, and is mostly single $\pi^0$ events. Sample D events have both Michel electrons and $\geq 2$ isolated clusters, and is mostly events with multiple pions. Details of these four samples are given in Ref. [25]. Sample A has $1.3 \times 10^6$ selected events with a predicted background of $0.4 \times 10^6$. Samples B–D contain $0.23 \times 10^6, 0.22 \times 10^6$, and $57\,000$ events.

For each bin of $p_T$ and $\Sigma T_p$, a joint fit to the above-listed four samples is used to determine scale factors applied to the signal sample (A) and to each of the backgrounds [single $\pi^+$ (B), single $\pi^0$ (C), and multipion (D)]. The fit minimizes a $\chi^2$ over the four scale factors using a singular value decomposition that drops singular values with condition number $\leq 10^{-3}$ to avoid numerical instability and fords negative scale factors for any component. The background-subtracted event rate is unfolded using an iterative technique [57] from the RooUnfold framework [58] that is regularized by the number of iterations. A regularization of 10 iterations was chosen by generating randomly fluctuated pseudodata samples with a number of different underlying physics models to ensure fidelity with different assumed data models. The statistical covariance matrix is scaled to account for the finite Monte Carlo statistics in the true-to-reconstructed migration matrix. The unfolded 3D distribution is then corrected for the predicted event loss from selection inefficiencies and detector effects. The average efficiency is between 40% and 75% over all event loss from selection inefficiencies and detector effects.

The unfolded $3\Sigma T_p$ distribution is then corrected for the predicted background of $3\Sigma T_p$, scaled to account for the finite Monte Carlo statistics in the true-to-reconstructed migration matrix. The statistical covariance matrix is regularized by the number of iterations. A regularization of 10 iterations was chosen by generating randomly fluctuated pseudodata samples with a number of different underlying physics models to ensure fidelity with different assumed data models. The statistical covariance matrix is scaled to account for the finite Monte Carlo statistics in the true-to-reconstructed migration matrix. The unfolded 3D distribution is then corrected for the predicted event loss from selection inefficiencies and detector effects. The average efficiency is between 40% and 75% over all event loss from selection inefficiencies and detector effects.

The average recoil energy in data falls $\sim 50$ MeV below the reference model at low $p_T \lesssim 0.5$ GeV/c, then rises to be comparable to the model $\sim 0.9$ GeV/c, and finally exceeds the model prediction in the highest $p_T$ bin. The abrupt change in the highest $p_T$ bins may be due to a cutoff in the Valencia multinucleon model that eliminates this process.

FIG. 2. Distribution of average $\Sigma T_p$ in $p_T - p_{\|}$ bins. Legends within each panel give the percentage contribution of the $p_T$ panel to the total integrated cross section for data and for the reference simulation. Statistical uncertainty is denoted by the colored box and the total uncertainty by the error bar. The “No Low Recoil Fit” prediction has no enhancement applied to the 2p2h model.
multinucleon and half low W pion production where the pion is absorbed in FSI. Therefore, it is likely that both processes are overpredicted. In this region, the multinucleon prediction is from the Valencia model with little effect from MINERvA tune v4.4.1. In pion production, such an effect may arise either from overprediction of the baryon-resonance pion production cross sections or from too-small suppression of primary pions due to FSI. A low $Q^2$ suppression of resonant pion production[18,61] or a reduction in the visible energy from these events [33], as Fig. 2 also suggests, would improve agreement with this data. Shifts in energy transfers have been observed in $(e,e')$ data for regions of low energy transfer [34,62–70]. An overprediction of pion FSI could arise from finite hadronic formation time [71], an effect not included in the reference simulation. However this background arises mostly from absorption of slow pions ($p_t < 0.3$ GeV/$c$), hence pion formation time is unlikely to account for the entire effect.

**Moderate $p_t$ and $\Sigma T_p$ just above the quasielastic peak.—** For $\Sigma T_p$ of 0.2 GeV and 0.15 < $p_t$ < 0.55 GeV/$c$, where the modifications of MINERvA tune v4.4.1 to multinucleon processes are large, the data and reference model would be in strong disagreement without these modifications. Figure 3 shows that the ratio of the data to the reference model dips near the peak of the tune, suggesting that the shape of the MINERvA tune v4.4.1 enhancement may not be accurate, either in rate or in fraction of events with a neutron in the final state. However, at

![Graph](image-url)
... $p_i > 0.55$ GeV/c, where the model predicts a smaller multinucleon contribution, the data mostly exceeds the reference prediction, suggesting that a significant enhancement to multinucleon processes at higher $p_i$ than in MINERvA tune v4.4.1 may be needed.

High $p_i$ and low $\Sigma T_p$: At $p_i > 0.55$ GeV/c and $\Sigma T_p < 50$ MeV, there is a significant overprediction relative to data. This region is dominated by true quasielastic events where the final-state proton undergoes FSI and leaves the nucleus as one or more energetic neutrons; this suggests that too much strength is given to FSI in this kinematic region.

Figure 4 presents the flux-averaged triple-differential cross section $d^3\sigma/dE_\mu dq_{0}^{(QE)}d\Sigma T_p$. Here as well, significant data versus reference model discrepancies are seen at low $q_{0}^{(QE)}$ for $\Sigma T_p$ beyond the peak of the quasielastic contribution. The previously noted discrepancy at low $\Sigma T_p$ and high $p_i$ corresponds to a predicted peak near zero $\Sigma T_p$ at high $q_{0}^{(QE)}$, which is absent from the data. The cross section in the quasielastic peak is underestimated, especially at higher $q_{0}^{(QE)}$, and modified form factors could improve this agreement [46,72–74]. This measurement directly probes the relationship between energy estimators used in oscillation experiments, and discrepancies with models suggest deficiencies in modeling those estimators.

In summary, a number of modeling shortfalls for neutrino-nucleus quasielastic-like scattering are identified by this measurement. These imply that relationships between the true neutrino energy of quasielastic-like events and experimental estimators, such as $E_{\text{available}} + E_\mu$ and $q_{0}^{(QE)} + E_\mu$, differ from those predicted by current neutrino generators. The triple-differential cross section presented in Fig. 4 can serve as a benchmark for neutrino-nucleus interaction simulations employed in ongoing and future neutrino oscillation experiments.

This document was prepared by members of the MINERvA Collaboration using the resources of the Fermi National Accelerator Laboratory (Fermilab), a U.S. Department of Energy, Office of Science, HEP User Facility. Fermilab is managed by Fermi Research Alliance, LLC (FRA), acting under Contract No. DE-AC02-07CH11359. These resources included support for the MINERvA construction project, and support for construction also was granted by the United States National Science Foundation under Grant No. PHY-0619727 and by the University of Rochester. Support for participating scientists was provided by NSF and DOE (USA); by CAPES and CNPq (Brazil); by Conicyt PIA ACT1413, and Fondecyt 3170845 and 11130133 (Chile); by CONCYTEC (Consejo Nacional de Ciencia, Tecnología e Innovación Tecnológica), DGI-PUCP (Dirección de Gestión de la Investigación—Pontificia Universidad Católica del Perú), and VRI-UNI (Vice-Rectorate for Research of National University of Engineering) (Peru); NCN Opus Grant No. 2016/21/B/ST2/01092 (Poland); by Science and Technology Facilities Council (UK); by EU Horizon 2020 Marie Skłodowska-Curie Action; by an Imperial College London President’s PhD Scholarship. D.R. gratefully acknowledges support from a Cottrell Postdoctoral Fellowship, Research Corporation for Scientific Advancement Grant No. 27467 and National Science Foundation Grant No. CHE2039044. We thank the MINOS Collaboration for use of its near detector data. Finally, we thank the staff of Fermilab for support of the beam line, the detector, and computing infrastructure.

---

*Present address: Iowa State University, Ames, Iowa 50011, USA.
††Present address: Iowa State University, Ames, Iowa 50011, USA.
‡Present address: Brookhaven National Laboratory, Upton, New York 11973-5000, USA.
∥Present address: Los Alamos National Laboratory, Los Alamos, New Mexico 87545, USA.
§Present address: Brookhaven National Laboratory, Upton, New York 11973-5000, USA.
∥∥Present address: Department of Physics and Astronomy, University of California at Davis, Davis, California 95616, USA.
**Present address: Brookhaven National Laboratory, Upton, New York 11973-5000, USA.
†††Present address: Department of Physics, University of Minnesota, Minneapolis, Minnesota 55455, USA.

et al. Measurements of Nuclear Effects and the \(\bar{\nu}_e + H \to \mu^+ + n\) Cross Section in MINERvA with Neutron Tagging, Ph.D. thesis, University of Rochester (2021).


59] See Supplemental Material at http://link.aps.org/supplemental/10.1103/PhysRevLett.129.021803 provides a separation of \(p_\parallel\) curves shown in Fig. 1 for a clear presentation of the uncertainty of the measurement and all \(p_\parallel\)(\(E_p\)) slices for Figs. 3 and 4 for a complete representation of the triple differential cross sections.


