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Evidence for CP Violation in Time-Integrated $D^0 \to h^- h^+$ Decay Rates

R. Aaij et al.*
(LHCb Collaboration)

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A search for time-integrated CP violation in $D^0 \to h^- h^+$ ($h = K, \pi$) decays is presented using 0.62 fb$^{-1}$ of data collected by LHCb in 2011. The flavor of the charm meson is determined by the charge of the slow pion in the $D^{\ast+} \to D^0 \pi^+$ and $D^{\ast-} \to D^0 \pi^-$ decay chains. The difference in CP asymmetry between $D^0 \to K^- K^+$ and $D^0 \to \pi^- \pi^+$, $\Delta A_{CP} = A_{CP}(K^- K^+) - A_{CP}(\pi^- \pi^+)$, is measured to be $[-0.82 \pm 0.21\,(stat) \pm 0.11\,(syst)]\%$. This differs from the hypothesis of CP reservation by 3.5 standard deviations.

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The charm sector is a promising place to probe for the effects of physics beyond the standard model (SM). There has been a resurgence of interest in the past few years since evidence for $D^0$ mixing was first seen [1,2]. Mixing is now well established [3] at a level which is consistent with, but at the upper end of, SM expectations [4]. By contrast, no evidence for CP violation in charm decays has yet been found.

The time-dependent CP asymmetry $A_{CP}(f; t)$ for $D^0$ decays to a CP eigenstate $f$ (with $f = \bar{f}$) is defined as

$$A_{CP}(f; t) = \frac{\Gamma(D^0(t) \to f) - \Gamma(D^0(t) \to \bar{f})}{\Gamma(D^0(t) \to f) + \Gamma(D^0(t) \to \bar{f})},$$

where $\Gamma$ is the decay rate for the process indicated. In general $A_{CP}(f; t)$ depends on $f$. For $f = K^- K^+$ and $f = \pi^- \pi^+$, $A_{CP}(f; t)$ can be expressed in terms of two contributions: a direct component associated with CP violation in the decay amplitudes, and an indirect component associated with CP violation in the mixing or in the interference between mixing and decay. In the limit of $U$-spin symmetry, the direct component is equal in magnitude and opposite in sign for $K^- K^+$ and $\pi^- \pi^+$, though the size of $U$-spin breaking effects remains to be quantified precisely [5].

The asymmetry $A_{CP}(f; t)$ may be written to first order as [10,11]

$$A_{CP}(f; t) = a_{CP}^{\dir}(f) + \frac{t}{\tau} a_{CP}^{\ind},$$

where $a_{CP}^{\dir}(f)$ is the direct CP asymmetry, $\tau$ is the $D^0$ lifetime, and $a_{CP}^{\ind}$ is the indirect CP asymmetry. To a good approximation this latter quantity is universal [5,12]. The time-integrated asymmetry measured by an experiment, $A_{CP}(f)$, depends upon the time acceptance of that experiment. It can be written as

$$A_{CP}(f) = a_{CP}^{\dir}(f) + \frac{\langle t \rangle}{\tau} a_{CP}^{\ind},$$

where $\langle t \rangle$ is the average decay time in the reconstructed sample. Denoting by $\Delta$ the differences between quantities for $D^0 \to K^- K^+$ and $D^0 \to \pi^- \pi^+$ it is then possible to write

$$\Delta A_{CP} \equiv A_{CP}(K^- K^+) - A_{CP}(\pi^- \pi^+) = [a_{CP}^{\dir}(K^- K^+) - a_{CP}^{\dir}(\pi^- \pi^+)] + \frac{\Delta(t)}{\tau} a_{CP}^{\ind}.\quad (4)$$

In the limit that $\Delta(t)$ vanishes, $\Delta A_{CP}$ is equal to the difference in the direct CP asymmetry between the two decays. However, if the time acceptance is different for the $K^- K^+$ and $\pi^- \pi^+$ final states, a contribution from indirect CP violation remains.

The most precise measurements to date of the time-integrated CP asymmetries in $D^0 \to K^- K^+$ and $D^0 \to \pi^- \pi^+$ were made by the CDF, BABAR, and Belle collaborations [10,13,14]. The Heavy Flavor Averaging Group (HFAG) has combined time-integrated and time-dependent measurements of CP asymmetries, taking account of the different decay time acceptances, to obtain world average values for the indirect CP asymmetry of $a_{CP}^{\ind} = (-0.03 \pm 0.23)\%$ and the difference in direct CP asymmetry between the final states of $\Delta a_{CP}^{\dir} = (-0.42 \pm 0.27)\%$ [3].

In this Letter, we present a measurement of the difference in time-integrated CP asymmetries between $D^0 \to K^- K^+$ and $D^0 \to \pi^- \pi^+$, performed with 0.62 fb$^{-1}$ of data collected at LHCb between March and June 2011. The flavor of the initial state ($D^0$ or $D^0$) is tagged by requiring a $D^{\ast+} \to D^0 \pi^+$ decay, with the flavor determined by the charge of the slow pion ($\pi^+$). The inclusion
of charge-conjugate modes is implied throughout, except in the definition of asymmetries.

The raw asymmetry for tagged $D^0$ decays to a final state $f$ is given by $A_{\text{raw}}(f)$, defined as

$$A_{\text{raw}}(f) = \frac{N(D^{++} \rightarrow D^0(f)\pi^+_s) - N(D^{+-} \rightarrow D^0(f)\pi^-_s)}{N(D^{++} \rightarrow D^0(f)\pi^+_s) + N(D^{+-} \rightarrow D^0(f)\pi^-_s)}.$$  

(5)

where $N(X)$ refers to the number of reconstructed events of decay $X$ after background subtraction.

To first order the raw asymmetries may be written as a sum of four components, due to physics and detector effects:

$$A_{\text{raw}}(f) = A_{\text{CP}}(f) + A_D(f) + A_D(\pi^+_s) + A_{\text{D}}(D^{++}).$$  

(6)

Here, $A_D(f)$ is the asymmetry in selecting the $D^0$ decay into the final state $f$, $A_D(\pi^+_s)$ is the asymmetry in selecting the slow pion from the $D^{++}$ decay chain, and $A_{\text{D}}(D^{++})$ is the production asymmetry for $D^{++}$ mesons. The asymmetries $A_D$ and $A_{\text{D}}$ are defined in the same fashion as $A_{\text{raw}}$.

The first-order expansion is valid since the individual asymmetries are small.

For a two-body decay of a spin-0 particle to a self-conjugate final state there can be no $D^0$ detection asymmetry, i.e., $A_D(K^-K^+) = A_D(\pi^- \pi^+) = 0$. Moreover, $A_D(\pi^+_s)$ and $A_{\text{D}}(D^{++})$ are independent of $f$ and thus in the first-order expansion of Eq. (5) those terms cancel in the difference $A_{\text{raw}}(K^-K^+) - A_{\text{raw}}(\pi^- \pi^+)$, resulting in

$$\Delta A_{\text{CP}} = A_{\text{raw}}(K^-K^+) - A_{\text{raw}}(\pi^- \pi^+).$$  

(7)

To minimize second-order effects that are related to the slightly different kinematic properties of the two decay modes and that do not cancel in $\Delta A_{\text{CP}}$, the analysis is performed in bins of the relevant kinematic variables, as discussed later.

The LHCb detector is a forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, and is described in detail in Ref. [15]. The Ring Imaging Cherenkov (RICH) detectors are of particular importance to this analysis, providing kaon-pion discrimination for the full range of track momenta used. The nominal downstream beam direction is aligned with the $+z$ axis, and the field direction in the LHCb dipole is such that charged particles are deflected in the horizontal ($xz$) plane. The field polarity was changed several times during data taking: about 60% of the data were taken with the down polarity and 40% with the other.

Selections are applied to provide samples of $D^{++} \rightarrow D^0(\pi^+_s)$ candidates, with $D^0 \rightarrow K^-K^+$ or $\pi^- \pi^+$. Events are required to pass both hardware and software trigger levels. A loose $D^0$ selection is applied in the final state of the software trigger, and in the offline analysis only candidates that are accepted by this trigger algorithm are considered. Both the trigger and offline selections impose a variety of requirements on kinematics and decay time to isolate the decays of interest, including requirements on the track fit quality, on the $D^0$ and $D^{++}$ vertex fit quality, on the transverse momentum ($p_T > 2$ GeV/$c$) and decay time ($\tau > 100$ m) of the $D^0$ candidate, on the angle between the $D^0$ momentum in the lab frame and its daughter momenta in the $D^0$ rest frame ($|\cos \theta| < 0.9$), that the $D^0$ trajectory points back to a primary vertex, and that the $D^0$ daughter tracks do not. In addition, the offline analysis exploits the capabilities of the RICH system to distinguish between pions and kaons when reconstructing the $D^0$ meson, with no tracks appearing as both pion and kaon candidates.

A fiducial region is implemented by imposing the requirement that the slow pion lies within the central part of the detector acceptance. This is necessary because the magnetic field bends pions of one charge to the left and those of the other charge to the right. For soft tracks at large angles in the $xz$ plane this implies that one charge is much more likely to remain within the 300 mrad horizontal detector acceptance, thus making $A_D(\pi^+_s)$ large. Although this asymmetry is formally independent of the $D^0$ decay mode, it breaks the assumption that the raw asymmetries are small and it carries a risk of second-order systematic effects if the ratio of efficiencies of $D^0 \rightarrow K^-K^+$ and $D^0 \rightarrow \pi^- \pi^+$ varies in the affected region. The fiducial requirements therefore exclude edge regions in the slow pion ($p_z, p$) plane. Similarly, a small region of phase space in which one charge of slow pion is more likely to be swept into the beam pipe region in the downstream tracking stations, and hence has reduced efficiency, is also excluded. After the implementation of the fiducial requirements about 70% of the events are retained.

The invariant mass spectra of selected $K^-K^+$ and $\pi^- \pi^+$ pairs are shown in Fig. 1. The half width at half maximum of the signal line shape is 8.6 MeV/$c^2$ for $K^-K^+$ and 11.2 MeV/$c^2$ for $\pi^- \pi^+$, where the difference is due to the kinematics of the decays and has no relevance for the subsequent analysis. The mass difference $\delta m$ spectra of selected candidates, where $\delta m = m_h(h^-h^+\pi^+_s) - m(h^-h^-\pi^-_s) - m(\pi^+_s)$ for $h = K$, $\pi$, are shown in Fig. 2. Candidates are required to lie inside a wide $\delta m$ window of 0–15 MeV/$c^2$, and in Fig. 2 and for all subsequent results candidates are in addition required to lie in a mass signal window of 1844–1884 MeV/$c^2$. The $D^{++}$ signal yields are approximately $1.44 \times 10^6$ in the $K^-K^+$ sample, and $0.38 \times 10^6$ in the $\pi^- \pi^+$ sample. Charm from $b$-hadron decays is strongly suppressed by the requirement that the $D^0$ originate from a primary vertex, and accounts for only 3% of the total yield. Of the events that contain at least one $D^{++}$ candidate, 12% contain more than one candidate; this is expected due to background soft pions from the primary vertex and all candidates are accepted. The background-subtracted average decay time of $D^0$ candidates passing the selection is measured for each final state, and the fractional difference $\Delta(\tau)/\tau$ is obtained.
Systematic uncertainties on this quantity are assigned for the uncertainty on the world average $D_0$ lifetime $\tau$ (0.04%), charm from $b$-hadron decays (0.18%), and the background-subtraction procedure (0.04%). Combining the systematic uncertainties in quadrature, we obtain $h_{\text{ti}}=9.83 \pm 0.0005$ ps, where the error is statistical only.

Fits are performed on the samples in order to determine $A_{\text{raw}}(K^-K^+)$ and $A_{\text{raw}}(\pi^-\pi^+)$. The production and detection asymmetries can vary with $p_T$ and pseudorapidity $\eta$, and so can the detection efficiency of the two different $D^0$ decays, in particular, through the effects of the particle identification requirements. The analysis is performed in 54 kinematic bins defined by the $p_T$ and $\eta$ of the $D^{*+}$ candidates, the momentum of the slow pion, and the sign of $p_x$ of the slow pion at the $D^{*+}$ vertex. The events are further partitioned in two ways. First, the data are divided between the two dipole magnet polarities. Second, the first 60% of data are processed separately from the remainder, with the division aligned with a break in data taking due to an LHC technical stop. In total, 216 statistically independent measurements are considered for each decay mode.

In each bin, one-dimensional unbinned maximum likelihood fits to the $\delta m$ spectra are performed. The signal is described as the sum of two Gaussian functions with a common mean $\mu$ but different widths $\sigma_i$, convolved with a function $B(\delta m; s) = \Theta(\delta m)\delta m^{i}$ taking account of the asymmetric shape of the measured $\delta m$ distribution. Here, $s \approx -0.975$ is a shape parameter fixed to the value determined from the global fits shown in Fig. 2, $\Theta$ is the Heaviside step function, and the convolution runs over $\delta m$. The background is described by an empirical function of the form $1 - e^{-(\delta m - \delta m_0)/\alpha}$, where $\delta m_0$ and $\alpha$ are free parameters describing the threshold and shape of the function, respectively. The $D^{*+}$ and $D^{*-}$ samples in a given bin are fitted simultaneously and share all shape parameters, except for a charge-dependent offset in the central value $\mu$ and an overall scale factor in the mass resolution. The raw asymmetry in the signal yields is extracted directly from this simultaneous fit. No fit parameters are shared between the 216 subsamples of data, nor between the $K^-K^+$ and $\pi^-\pi^+$ final states.

The fits do not distinguish between the signal and backgrounds that peak in $\delta m$. Such backgrounds can arise from $D^{*+}$ decays in which the correct slow pion is found but the

![FIG. 1 (color online). Fits to the (a) $m(K^-K^+)$ and (b) $m(\pi^-\pi^+)$ spectra of $D^{*+}$ candidates passing the selection and satisfying $0 < \delta m < 15$ MeV/$c^2$. The dashed line corresponds to the background component in the fit, and the vertical lines indicate the signal window of 1844–1884 MeV/$c^2$.](image1)

![FIG. 2 (color online). Fits to the $\delta m$ spectra, where the $D^0$ is reconstructed in the final states (a) $K^-K^+$ and (b) $\pi^-\pi^+$, with mass lying in the window of 1844–1884 MeV/$c^2$. The dashed line corresponds to the background component in the fit.](image2)
D⁰ is partially misreconstructed. These backgrounds are suppressed by the use of tight particle identification requirements and a narrow D⁰ mass window. From studies of the D⁰ mass sidebands (1820–1840 and 1890–1910 MeV/c²), this contamination is found to be approximately 1% of the signal yield and to have small raw asymmetry (consistent with zero asymmetry difference between the K⁻K⁺ and π⁻π⁺ final states). Its effect on the measurement is estimated in an ensemble of simulated experiments and found to be negligible; a systematic uncertainty is assigned below based on the statistical precision of the estimate.

A value of ΔACP is determined in each measurement bin as the difference between A_avg(K⁻K⁺) and A_avg(π⁻π⁺). Testing these 216 measurements for mutual consistency, we obtain χ²/ndf = 211/215 (χ² probability of 56%). A weighted average is performed to yield the result ΔACP = (−0.82 ± 0.21)%), where the uncertainty is statistical only.

Numerous robustness checks are made. The value of ΔACP is studied as a function of the time at which the data were taken (Fig. 3) and found to be consistent with a constant value (χ² probability of 57%). The measurement is repeated with progressively more restrictive RICH particle identification requirements, finding values of (−0.88 ± 0.26)% and (−1.03 ± 0.31)%; both of these values are consistent with the baseline result when correlations are taken into account. Table I lists ΔACP for eight disjoint subsamples of data split according to magnet polarities, the sign of px of the slow pion, and whether the data were taken before or after the technical stop. The χ² probability for consistency among the subsamples is 45%. The significances of the differences between data taken before and after the technical stop, between the magnet polarities, and between px > 0 and px < 0 are 0.4, 0.6, and 0.7 standard deviations, respectively. Other checks include applying electron and muon vetoes to the slow pion and to the D⁰ daughters, use of different kinematic binnings, validation of the size of the statistical uncertainties with Monte Carlo pseudoexperiments, tightening of kinematic requirements, testing for variation of the result with the multiplicity of tracks and of primary vertices in the event, use of other signal and background parameterizations in the fit, and imposing a full set of common shape parameters between D⁺⁺ and D⁻⁻ candidates. Potential biases due to the inclusive hardware trigger selection are investigated with the subsample of data in which one of the signal final-state tracks is directly responsible for the hardware trigger decision. In all cases good stability is observed. For several of these checks, a reduced number of kinematic bins are used for simplicity. No systematic dependence of ΔACP is observed with respect to the kinematic variables.

Systematic uncertainties are assigned by loosening the fiducial requirement on the slow pion, assessing the effect of potential peaking backgrounds in Monte Carlo pseudoexperiments, repeating the analysis with the asymmetry extracted through sideband subtraction in δm instead of a fit, removing all candidates but one (chosen at random) in events with multiple candidates, and comparing with the result obtained without kinematic binning. In each case the

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### Table I

<table>
<thead>
<tr>
<th>Subsample</th>
<th>ΔACP [%]</th>
<th>χ²/ndf</th>
</tr>
</thead>
<tbody>
<tr>
<td>Pre-TS, up, left</td>
<td>−1.22 ± 0.59</td>
<td>13/26(98%)</td>
</tr>
<tr>
<td>Pre-TS, up, right</td>
<td>−1.43 ± 0.59</td>
<td>27/26(39%)</td>
</tr>
<tr>
<td>Pre-TS, down, left</td>
<td>−0.59 ± 0.52</td>
<td>19/26(84%)</td>
</tr>
<tr>
<td>Pre-TS, down, right</td>
<td>−0.51 ± 0.52</td>
<td>29/26(30%)</td>
</tr>
<tr>
<td>Post-TS, up, left</td>
<td>−0.79 ± 0.90</td>
<td>26/26(44%)</td>
</tr>
<tr>
<td>Post-TS, up, right</td>
<td>+0.42 ± 0.93</td>
<td>21/26(77%)</td>
</tr>
<tr>
<td>Post-TS, down, left</td>
<td>−0.24 ± 0.56</td>
<td>34/26(15%)</td>
</tr>
<tr>
<td>Post-TS, down, right</td>
<td>−1.59 ± 0.57</td>
<td>35/26(12%)</td>
</tr>
<tr>
<td>All data</td>
<td>−0.82 ± 0.21</td>
<td>211/215(56%)</td>
</tr>
</tbody>
</table>

---

### Table II

Summary of absolute systematic uncertainties for ΔACP.

<table>
<thead>
<tr>
<th>Source</th>
<th>Uncertainty</th>
</tr>
</thead>
<tbody>
<tr>
<td>Fiducial requirement</td>
<td>0.01%</td>
</tr>
<tr>
<td>Peaking background asymmetry</td>
<td>0.04%</td>
</tr>
<tr>
<td>Fit procedure</td>
<td>0.08%</td>
</tr>
<tr>
<td>Multiple candidates</td>
<td>0.06%</td>
</tr>
<tr>
<td>Kinematic binning</td>
<td>0.02%</td>
</tr>
<tr>
<td>Total</td>
<td>0.11%</td>
</tr>
</tbody>
</table>
full value of the change in result is taken as the systematic uncertainty. These uncertainties are listed in Table II. The sum in quadrature of the statistical and systematic uncertainties is 0.22(stat) ± 0.19(sys). The contribution from indirect CP violation is suppressed and ΔACP is primarily sensitive to direct CP violation. Dividing the central value by the sum in quadrature of the statistical and systematic uncertainties, the significance of the measured deviation from zero is 3.5σ. This is the first evidence for CP violation in the charm sector. To establish whether this result is consistent with the SM will require the analysis of more data, as well as improved theoretical understanding.

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STFC Rutherford Appleton Laboratory, Didcot, United Kingdom

School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom

School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom

Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom

Imperial College London, London, United Kingdom

School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom

Department of Physics, University of Oxford, Oxford, United Kingdom

Syracuse University, Syracuse, New York, USA

CC-IN2P3, CNRS/IN2P3, Lyon-Villeurbanne, France

Pontificia Universidade Católica do Rio de Janeiro (PUC-Rio), Rio de Janeiro, Brazil

University of Birmingham, Birmingham, United Kingdom

Physikalisches Institut, Universität Rostock, Rostock, Germany

aAlso at P.N. Lebedev Physical Institute, Russian Academy of Science (LPI RAS), Moscow, Russia.
bAlso at Università di Bari, Bari, Italy.
cAlso at Università di Bologna, Bologna, Italy.
dAlso at Università di Cagliari, Cagliari, Italy.
eAlso at Università di Ferrara, Ferrara, Italy.
fAlso at Università di Firenze, Firenze, Italy.
gAlso at Università di Urbino, Urbino, Italy.
hAlso at Università di Modena e Reggio Emilia, Modena, Italy.
iAlso at Università di Genova, Genova, Italy.
jAlso at Università di Milano Bicocca, Milano, Italy.
kAlso at Università di Roma Tor Vergata, Roma, Italy.
lAlso at Università di Roma La Sapienza, Roma, Italy.
mAlso at Università della Basilicata, Potenza, Italy.
nAlso at LIFAELS, La Salle, Universitat Ramon Llull, Barcelona, Spain.
oAlso at Hanoi University of Science, Hanoi, Vietnam.