Measurements of the branching fractions of $B^+ \to p\bar{p}K^+$ decays

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Abstract The branching fractions of the decay $B^+ \to p\bar{p}K^+$ for different intermediate states are measured using data, corresponding to an integrated luminosity of $1.0 \text{ fb}^{-1}$, collected by the LHCb experiment. The total branching fraction, its charmless component ($M_{p\bar{p}} < 2.85 \text{ GeV/c}^2$) and the branching fractions via the resonant $c\bar{c}$ states $\eta_c(1S)$ and $\psi(2S)$ relative to the decay via a $J/\psi$ intermediate state are

$$\frac{\mathcal{B}(B^+ \to p\bar{p}K^+)}{\mathcal{B}(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 4.91 \pm 0.19 \text{ (stat)} \pm 0.14 \text{ (syst)},$$

$$\frac{\mathcal{B}(B^+ \to p\bar{p}K^+)}{\mathcal{B}(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 2.02 \pm 0.10 \text{ (stat)} \pm 0.08 \text{ (syst)},$$

$$\frac{\mathcal{B}(B^+ \to \eta_c(1S)K^+ \to p\bar{p}K^+)}{\mathcal{B}(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 0.578 \pm 0.035 \text{ (stat)} \pm 0.027 \text{ (syst)},$$

$$\frac{\mathcal{B}(B^+ \to \psi(2S)K^+ \to p\bar{p}K^+)}{\mathcal{B}(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 0.080 \pm 0.012 \text{ (stat)} \pm 0.009 \text{ (syst)}.$$

Upper limits on the $B^+$ branching fractions into the $\eta_c(2S)$ meson and into the charmonium-like states $X(3872)$ and $X(3915)$ are also obtained.

1 Introduction

The $B^+ \to p\bar{p}K^+$ decay\(^{1}\) offers a clean environment to study $c\bar{c}$ states and charmonium-like mesons that decay to $p\bar{p}$ and excited $\Lambda$ baryons that decay to $p\bar{p}K^+$, and to search for glueballs or exotic states. The presence of $p\bar{p}$ in the final state allows intermediate states of any quantum numbers to be studied and the existence of the charged kaon in the final state significantly enhances the signal to background ratio in the selection procedure. Measurements of intermediate charmonium-like states, such as the $X(3872)$, are important to clarify their nature \cite{1,2} and to determine their partial width to $p\bar{p}$, which is crucial to predict the production rate of these states in dedicated experiments \cite{3}. BaBar and Belle have previously measured the $B^+ \to p\bar{p}K^+$ branching fraction, including contributions from the $J/\psi$ and $\eta_c(1S)$ intermediate states \cite{4,5}. The data sample, corresponding to an integrated luminosity of $1.0 \text{ fb}^{-1}$, collected by LHCb at $\sqrt{s} = 7 \text{ TeV}$ allows the study of substructures in the $B^+ \to p\bar{p}K^+$ decays with a sample ten times larger than those available at previous experiments.

In this paper we report measurements of the ratios of branching fractions

$$\mathcal{R}(\text{mode}) = \frac{\mathcal{B}(B^+ \to \text{mode} \to p\bar{p}K^+)}{\mathcal{B}(B^+ \to J/\psi K^+ \to p\bar{p}K^+)},$$

where “mode” corresponds to the intermediate $\eta_c(1S)$, $\psi(2S)$, $\eta_c(2S)$, $\chi_{c0}(1P)$, $h_c(1P)$, $X(3872)$ or $X(3915)$ states, together with a kaon.

2 Detector and software

The LHCb detector \cite{6} is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing $b$ or $c$ quarks. The detector includes a high precision tracking system consisting of a silicon-strip vertex detector surrounding the $pp$ interaction region, a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift-tubes placed downstream. The combined tracking system has momentum ($p$) resolution $\Delta p/p$ that varies from 0.4% at 5 GeV/c to 0.6% at 100 GeV/c, and impact parameter resolution of 20 $\mu$m for tracks with high transverse...
momentum ($p_T$). Charged hadrons are identified using two ring-imaging Cherenkov (RICH) detectors. Photon, electron and hadron candidates are identified by a calorimeter system consisting of scintillating-pad and pre-shower detectors, an electromagnetic calorimeter and a hadronic calorimeter. Muons are identified by a system composed of alternating layers of iron and multiwire proportional chambers.

The trigger [7] consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage where candidates are fully reconstructed. The hardware trigger selects hadrons with high transverse energy in the calorimeter. The software trigger requires a two-, three- or four-track secondary vertex with a high $p_T$ sum of the tracks and a significant displacement from the primary vertex in the event. Particle identification (PID) responses of the electrons and hadrons are implemented to evaluate the ratio of the efficiencies for each considered channel with respect to the primary interaction greater than 16. The variables used by the BDT to discriminate between signal and background candidates are: the $p_T$ of each reconstructed track; the sum of the daughters’ $p_T$; the sum of the IP $\chi^2$ of the three daughter tracks with respect to the primary vertex; the IP of the daughter with the highest $p_T$, with respect to the primary vertex; the number of daughters with $p_T > 900$ GeV/c; the maximum distance of closest approach between any two of the $B^+$ daughter particles; the IP of the $B^+$ candidate with respect to the primary vertex; the distance between primary and secondary vertices; the $\theta_0$ angle; the $\chi^2$/ndf of the secondary vertex; a pointing variable defined as $\frac{P \sin \theta}{\sum_i p_{T,i}}$, where $P$ is the total momentum of the three-particle final state, $\theta$ is the angle between the direction of the sum of the daughter’s momentum and the direction of the flight distance of the $B^+$ and $\sum_i p_{T,i}$ is the sum of the transverse momenta of the daughters; and the log likelihood difference for each daughter between the assumed PID hypothesis and the pion hypothesis. The selection criterion on the BDT response (Fig. 1) is chosen in order to have a signal to background ratio of the order of unity. This corresponds to a BDT response value of −0.11. The efficiency of the BDT selection is greater than 92 % with a background rejection greater than 86 %.

4 Signal yield determination

The signal yield is determined from an unbinned extended maximum likelihood fit to the invariant mass of selected $B^+ \to p\bar{p}K^+$ candidates, shown in Fig. 2(a). The signal component is parametrized as the sum of two Gaussian functions with the same mean and different widths. The background component is parametrized as a linear function. The signal yield of the charmless component is determined by
performing the same fit described above to the sample of $B^+ \rightarrow p \bar{p} K^+$ candidates with $M_{p\bar{p}} < 2.85 \text{ GeV}/c^2$, shown in Fig. 2(b). The $B^+$ mass and widths, evaluated with the invariant mass fits to all of the $B^+ \rightarrow p \bar{p} K^+$ candidates, are compatible with the values obtained for the charmless component.

The signal yields for the charmonium contributions, $B^+ \rightarrow (cc) K^+ \rightarrow p \bar{p} K^+$, are determined by fitting the $p \bar{p}$ invariant mass distribution of $B^+ \rightarrow p \bar{p} K^+$ candidates within the $B^+$ mass signal window, $|M_{p\bar{p}K^+} - M_{B^+}| < 50 \text{ MeV}/c^2$. Simulations show that no narrow structures are induced in the $p \bar{p}$ spectrum as kinematic reflections of possible $B^+ \rightarrow p \Lambda \rightarrow p \bar{p} K^+$ intermediate states.

An unbinned extended maximum likelihood fit to the $p \bar{p}$ invariant mass distribution, shown in Fig. 3, is performed over the mass range 2400–4500 MeV/$c^2$. The signal components of the narrow resonances $J/\psi$, $\psi(2S)$, $h_c(1P)$, and $X(3872)$, whose natural widths are much smaller than the $p \bar{p}$ invariant mass resolution, are parametrized by Gaussian functions. The signal components for the $\eta_c(1S)$, $\chi_{c0}(1P)$, $\eta_c(2S)$, and $X(3915)$ are parametrized by Voigtian functions. Since the $p \bar{p}$ invariant mass resolution is approximately constant in the explored range, the resolution parameters for all resonances, except the $\psi(2S)$, are fixed to the $J/\psi$ value ($\sigma_{J/\psi} = 8.9 \pm 0.2 \text{ MeV}/c^2$). The background shape is parametrized as $f(M) = e^{c_1 M + c_2 M^2}$ where $c_1$ and $c_2$ are fit parameters. The $J/\psi$ and $\psi(2S)$ resolution parameters, the mass values of the $\eta_c(1S)$, $J/\psi$ and $\psi(2S)$ states, and the $\eta_c(1S)$ natural width are left free in the fit. The masses and widths for the other signal components are fixed to the corresponding world averages [16]. The $p \bar{p}$ invariant mass resolution, determined by the fit to the $\psi(2S)$ is $\sigma_{\psi(2S)} = 7.9 \pm 1.7 \text{ MeV}/c^2$.

The fit result is shown in Fig. 3. Figures 4 and 5 show the details of the fit result in the regions around the $\eta_c(1S)$ and $J/\psi$, $\eta_c(2S)$ and $\psi(2S)$, $\chi_{c0}(1P)$ and $h_c(1P)$, and $X(3872)$ and $X(3915)$ resonances. Any bias introduced by the inaccurate description of the tails of the $\eta_c(1S)$, $J/\psi$ and $\psi(2S)$ resonances is taken into account in the systematic uncertainty evaluation.

The contribution of $c \bar{c} \rightarrow p \bar{p} f$ from processes other than $B^+ \rightarrow p \bar{p} K^+$ decays, denoted as “non-signal”, is estimated from a fit to the $p \bar{p}$ mass in the $B^+$ mass sidebands 5130–5180 and 5380–5430 MeV/$c^2$. Except for the $J/\psi$ mode, no evidence of a non-signal contribution is found. The non-signal contribution to the $J/\psi$ signal yield in the $B^+$ mass window is $43 \pm 11$ candidates and is subtracted from the number of $J/\psi$ signal candidates.

The signal yields, corrected for the non-signal contribution, are reported in Table 1. For the intermediate charmonium states $\eta_c(2S)$, $\chi_{c0}(1P)$, $h_c(1P)$, $X(3872)$ and

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2A Voigtian function is the convolution of a Breit-Wigner function with a Gaussian distribution.
Fig. 3 Invariant mass distribution of the $p\bar{p}$ system for $B^+ \rightarrow p\bar{p}K^+$ candidates within the $B^+$ mass signal window, $|M(p\bar{p}K^+) - M_{B^+}| < 50$ MeV/$c^2$. The dotted lines represent the Gaussian and Voigtian functions (red) and the dashed line the smooth function (green) used to parametrize the signal and the background, respectively. The bottom plot shows the pulls (Color figure online).

Fig. 4 Invariant mass distribution of the $p\bar{p}$ system in the regions around (a) the $\eta_c(1S)$ and $J/\psi$ and (b) the $\eta_c(2S)$ and $\psi(2S)$ states. The dotted lines represent the Gaussian and the Voigtian functions (red) and the dashed line the smooth function (green) used to parametrize the signal and the background, respectively. The two plots below the mass distribution show the pulls (Color figure online).

Fig. 5 Invariant mass distribution of the $p\bar{p}$ system in the regions around (a) the $X(3872)$ and $h_c$, and (b) the $X(3872)$ and $X(3915)$ states. The dotted lines represent the Gaussian and Voigtian functions (red) and the dashed line the smooth function (green) used to parametrize the signal and the background, respectively. The two plots below the mass distribution show the pulls (Color figure online).

Table 1 Signal yields for the different channels and corresponding 95 % CL upper limits for modes with less than 3 $\sigma$ statistical significance. For the $J/\psi$ mode, the non-signal yield is subtracted. Uncertainties are statistical only.

<table>
<thead>
<tr>
<th>$B^+$ decay mode</th>
<th>Signal yield</th>
<th>Upper limit (95 % CL)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$p\bar{p}K^+$ [total]</td>
<td>$6951 \pm 176$</td>
<td></td>
</tr>
<tr>
<td>$p\bar{p}K^+$ [$M_{p\bar{p}} &lt; 2.85$ GeV/$c^2$]</td>
<td>$3238 \pm 122$</td>
<td></td>
</tr>
<tr>
<td>$J/\psi K^+$</td>
<td>$1458 \pm 42$</td>
<td></td>
</tr>
<tr>
<td>$\eta_c(1S)K^+$</td>
<td>$856 \pm 46$</td>
<td></td>
</tr>
<tr>
<td>$\psi(2S)K^+$</td>
<td>$107 \pm 16$</td>
<td></td>
</tr>
<tr>
<td>$\eta_c(2S)K^+$</td>
<td>$39 \pm 15$</td>
<td>$&lt;65.4$</td>
</tr>
<tr>
<td>$X(1P)K^+$</td>
<td>$15 \pm 13$</td>
<td>$&lt;38.1$</td>
</tr>
<tr>
<td>$h_c(1P)K^+$</td>
<td>$21 \pm 11$</td>
<td>$&lt;40.2$</td>
</tr>
<tr>
<td>$X(3872)K^+$</td>
<td>$-9 \pm 8$</td>
<td>$&lt;10.3$</td>
</tr>
<tr>
<td>$X(3915)K^+$</td>
<td>$13 \pm 17$</td>
<td>$&lt;42.1$</td>
</tr>
</tbody>
</table>

are determined from the likelihood profile integrating over the nuisance parameters. Since for the $X(3872)$ the fitted signal yield is negative, the upper limit has been calculated.
integrating the likelihood only in the physical region of a signal yield greater than zero.

5 Efficiency determination

The ratio of branching fractions is calculated using

$$R(\text{mode}) = \frac{B(B^+ \to \text{mode} \to p\bar{p}K^+)}{B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = \frac{N_{\text{mode}}}{N_{J/\psi}} \times \frac{\varepsilon_{J/\psi}}{\varepsilon_{\text{mode}}}, \quad (2)$$

where $N_{\text{mode}}$ and $N_{J/\psi}$ are the signal yields for the given mode and the reference mode, $B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)$, and $\varepsilon_{\text{mode}}/\varepsilon_{J/\psi}$ is the corresponding ratio of efficiencies. The efficiency is the product of the reconstruction, trigger, and selection efficiencies, and is estimated using simulated data samples.

Since the track multiplicity distribution for simulated events differs from that observed in data, simulated candidates are assigned a weight so that the weighted distribution reproduces the observed multiplicity distribution. The distributions of $\Delta \ln \mathcal{L}_{K\pi}$ and $\Delta \ln \mathcal{L}_{p\pi}$ for kaons and protons in data are obtained in bins of momentum, pseudorapidity and number of tracks from control samples of $D^{*-} \to D^0 (\to K^-\pi^+\pi^-)\pi^+$ decays for kaons and $\Lambda \to p\pi^-$ decays for protons, which are then used on a track-by-track basis to correct the simulation. The efficiency as a function of $M_{p\bar{p}}$ is shown in Fig. 6. A linear fit to the efficiency distribution is performed and the efficiency ratios are determined based on the fit result.

6 Systematic uncertainties

The measurements of the relative branching fractions depend on the ratios of signal yields and efficiencies with respect to the reference mode. Since the final state is the same in all cases, most of the systematic uncertainties cancel. The systematic uncertainty on the efficiency ratio, in each region of $p\bar{p}$ invariant mass, is determined from the difference between the efficiency ratios calculated using the solid fitted line and the dashed point-by-point interpolation shown in Fig. 6. The uncertainty associated with the evaluation of the $B^+$ signal yield has been determined by varying the fit range by $\pm 30$ MeV/$c^2$, using a single Gaussian instead of a double Gaussian function to model the signal PDF, and using an exponential function to model the background. For each charmonium resonance the systematic uncertainty on the signal yield has been investigated by varying the $B$ mass signal window by $\pm 10$ MeV/$c^2$, the signal and background shape parametrization and the subtraction of the $c\bar{c}$ contribution from the continuum. The systematic uncertainty associated with the parametrization of the signal tails of the $J/\psi, \eta_c(1S)$ and $\psi(2S)$ resonances is taken into account by taking the difference between the number of candidates in the observed distribution and the number of candidates calculated from the integral of the fit function in the range $-6\sigma$ to $-2.5\sigma$. The systematic uncertainty associated with the selection procedure is estimated by changing the value of the BDT selection to $-0.03$, which retains 85% of the signal with a 30% background, and is found to be negligible. The contributions to the systematic uncertainties from the different sources are listed in Table 2. The total systematic uncertainty is determined by adding the individual contributions in quadrature.

7 Results

The results are summarized in Table 3 and the values of the product of branching fractions derived from our measurement using the world average values $B(B^+ \to J/\psi K^+) = (1.013 \pm 0.034) \times 10^{-3}$ and $B(J/\psi \to p\bar{p}) = (2.17 \pm 0.07) \times 10^{-3}$ [16] are listed in Table 4. The branching fractions obtained are compatible with the world average values [16]. The upper limit on $B(B^+ \to \chi_{c0}(1P)K^+ \to p\bar{p}K^+)$ is compatible with the world average $B(B^+ \to \chi_{c0}(1P)K^+ \times B(\chi_{c0}(1P) \to p\bar{p}) = (0.030 \pm 0.004) \times 10^{-6}$ [16]. We combine our upper limit for $X(3872)$ with the known value for $B(B^+ \to X(3872)K^+) \times B(X(3872) \to J/\psi\pi^+\pi^-) = (8.6 \pm 0.8) \times 10^{-6}$ [16] to obtain the limit

$$\frac{B(X(3872) \to p\bar{p})}{B(X(3872) \to J/\psi\pi^+\pi^-)} < 2.0 \times 10^{-3}.$$
Table 2  Relative systematic uncertainties (in %) on the relative branching fractions from different sources. The total systematic uncertainty is determined by adding the individual contributions in quadrature.

<table>
<thead>
<tr>
<th>Source</th>
<th>$\mathcal{R}$(total)</th>
<th>$\mathcal{R}(M_{p\bar{p}} &lt; 2.85\text{ GeV/c}^2)$</th>
<th>$\mathcal{R}(\eta_c(1S))$</th>
<th>$\mathcal{R}(\psi(2S))$</th>
</tr>
</thead>
<tbody>
<tr>
<td>Efficiency ratio</td>
<td>0.21</td>
<td>0.5</td>
<td>3.3</td>
<td>4.8</td>
</tr>
<tr>
<td>$B^+$ mass fit range</td>
<td>0.16</td>
<td>0.5</td>
<td>–</td>
<td>–</td>
</tr>
<tr>
<td>Sig. and Bkg. shape</td>
<td>2.5</td>
<td>3.6</td>
<td>1.8</td>
<td>6.5</td>
</tr>
<tr>
<td>$B^+$ mass window</td>
<td>0.6</td>
<td>0.6</td>
<td>0.9</td>
<td>3.8</td>
</tr>
<tr>
<td>Non-signal component</td>
<td>–</td>
<td>–</td>
<td>0.4</td>
<td>5.1</td>
</tr>
<tr>
<td>Signal tail param.</td>
<td>1.0</td>
<td>1.0</td>
<td>1.2</td>
<td>4.3</td>
</tr>
<tr>
<td>Total</td>
<td>2.8</td>
<td>3.8</td>
<td>4.1</td>
<td>11.3</td>
</tr>
</tbody>
</table>

Table 3  Signal yields, efficiency ratios, ratios of branching fractions and corresponding upper limits.

<table>
<thead>
<tr>
<th>$B^+ \to (\text{mode}) \to p\bar{p}K^+$</th>
<th>Yield $\pm$ stat $\pm$ syst</th>
<th>$\epsilon_{\text{mode}}/\epsilon_{J/\psi}$ $\pm$ syst</th>
<th>$\mathcal{R}$(mode) $\pm$ stat $\pm$ syst</th>
<th>Upper Limit 95 % CL</th>
</tr>
</thead>
<tbody>
<tr>
<td>$J/\psi K^+$</td>
<td>1458 $\pm$ 42 $\pm$ 24</td>
<td>–</td>
<td>1</td>
<td>–</td>
</tr>
<tr>
<td>total</td>
<td>6951 $\pm$ 176 $\pm$ 171</td>
<td>0.970 $\pm$ 0.002 $\pm$ 0.000</td>
<td>4.91 $\pm$ 0.19 $\pm$ 0.14</td>
<td>–</td>
</tr>
<tr>
<td>$M_{p\bar{p}} &lt; 2.85\text{ GeV/c}^2$</td>
<td>3238 $\pm$ 122 $\pm$ 121</td>
<td>1.097 $\pm$ 0.006 $\pm$ 0.000</td>
<td>2.02 $\pm$ 0.10 $\pm$ 0.08</td>
<td>–</td>
</tr>
<tr>
<td>$\eta_c(1S)K^+$</td>
<td>856 $\pm$ 46 $\pm$ 19</td>
<td>1.016 $\pm$ 0.034 $\pm$ 0.009</td>
<td>0.578 $\pm$ 0.035 $\pm$ 0.026</td>
<td>–</td>
</tr>
<tr>
<td>$\psi(2S)K^+$</td>
<td>107 $\pm$ 16 $\pm$ 13</td>
<td>0.921 $\pm$ 0.044 $\pm$ 0.009</td>
<td>0.080 $\pm$ 0.012 $\pm$ 0.009</td>
<td>–</td>
</tr>
<tr>
<td>$\eta_c(2S)K^+$</td>
<td>39 $\pm$ 15 $\pm$ 5</td>
<td>0.927 $\pm$ 0.041 $\pm$ 0.009</td>
<td>0.029 $\pm$ 0.011 $\pm$ 0.004</td>
<td>&lt;0.048</td>
</tr>
<tr>
<td>$\chi_{c0}(1P)K^+$</td>
<td>15 $\pm$ 13 $\pm$ 4</td>
<td>0.957 $\pm$ 0.024 $\pm$ 0.009</td>
<td>0.011 $\pm$ 0.009 $\pm$ 0.003</td>
<td>&lt;0.028</td>
</tr>
<tr>
<td>$h_c(1P)K^+$</td>
<td>21 $\pm$ 11 $\pm$ 5</td>
<td>0.943 $\pm$ 0.032 $\pm$ 0.009</td>
<td>0.015 $\pm$ 0.008 $\pm$ 0.004</td>
<td>&lt;0.029</td>
</tr>
<tr>
<td>$X(3872)K^+$</td>
<td>$-9\pm8\pm2$</td>
<td>0.896 $\pm$ 0.058 $\pm$ 0.007</td>
<td>$-0.007\pm0.006\pm0.002$</td>
<td>&lt;0.008</td>
</tr>
<tr>
<td>$X(3915)K^+$</td>
<td>$13\pm17\pm5$</td>
<td>0.890 $\pm$ 0.062 $\pm$ 0.004</td>
<td>0.010 $\pm$ 0.013 $\pm$ 0.002</td>
<td>&lt;0.032</td>
</tr>
</tbody>
</table>

Table 4  Branching fractions for $B^+ \to (\text{mode}) \to p\bar{p}K^+$ derived using the world average value of the $B(B^+ \to J/\psi K^+)$ and $B(J/\psi \to p\bar{p})$ branching fractions [16]. For the charmonium modes we compare our values to the product of the independently measured branching fractions. The first uncertainties are statistical, the second systematic in the present measurement, and the third systematic from the uncertainty on the $J/\psi$ branching fraction.

<table>
<thead>
<tr>
<th>$B^+$ decay mode</th>
<th>$B(B^+ \to (\text{mode}) \to p\bar{p}K^+) \times 10^6$</th>
<th>UL (95 % CL) $\times 10^6$</th>
<th>Previous measurements</th>
</tr>
</thead>
<tbody>
<tr>
<td>total</td>
<td>10.81 $\pm$ 0.42 $\pm$ 0.30 $\pm$ 0.49</td>
<td>10.76$^{+0.36}_{-0.33}$ $\pm$ 0.70</td>
<td></td>
</tr>
<tr>
<td>$M_{p\bar{p}} &lt; 2.85\text{ GeV/c}^2$</td>
<td>4.46 $\pm$ 0.21 $\pm$ 0.18 $\pm$ 0.20</td>
<td>5.12 $\pm$ 0.31</td>
<td></td>
</tr>
<tr>
<td>$\eta_c(1S)K^+$</td>
<td>1.27 $\pm$ 0.08 $\pm$ 0.05 $\pm$ 0.06</td>
<td>1.54 $\pm$ 0.16</td>
<td></td>
</tr>
<tr>
<td>$\psi(2S)K^+$</td>
<td>0.175 $\pm$ 0.027 $\pm$ 0.020 $\pm$ 0.008</td>
<td>0.176 $\pm$ 0.012</td>
<td></td>
</tr>
<tr>
<td>$\eta_c(2S)K^+$</td>
<td>0.063 $\pm$ 0.025 $\pm$ 0.009 $\pm$ 0.003</td>
<td>&lt;0.106</td>
<td></td>
</tr>
<tr>
<td>$\chi_{c0}(1P)K^+$</td>
<td>0.024 $\pm$ 0.021 $\pm$ 0.006 $\pm$ 0.001</td>
<td>&lt;0.062</td>
<td>0.030 $\pm$ 0.004</td>
</tr>
<tr>
<td>$h_c(1P)K^+$</td>
<td>0.034 $\pm$ 0.018 $\pm$ 0.008 $\pm$ 0.002</td>
<td>&lt;0.064</td>
<td></td>
</tr>
<tr>
<td>$X(3872)K^+$</td>
<td>$-0.015\pm0.013\pm0.003\pm0.001$</td>
<td>&lt;0.017</td>
<td></td>
</tr>
<tr>
<td>$X(3915)K^+$</td>
<td>0.022 $\pm$ 0.029 $\pm$ 0.004 $\pm$ 0.001</td>
<td>&lt;0.071</td>
<td></td>
</tr>
</tbody>
</table>
This limit challenges some of the predictions for the molecular interpretations of the $X(3872)$ state and is approaching the range of predictions for a conventional $\chi_c(2P)$ state [17, 18]. Using our result and the $\eta_c(2S)$ branching fraction $B(\eta_c(2S) \to K\bar{K}\pi) = (3.4^{+2.3}_{-1.6}) \times 10^{-6}$ [16], a limit of

$$\frac{B(\eta_c(2S) \to p\bar{p})}{B(\eta_c(2S) \to K\bar{K}\pi)} < 3.1 \times 10^{-2}$$

is obtained.

8 Summary

Based on a sample of $6951 \pm 176$ $B^+ \to p\bar{p}K^+$ decays reconstructed in a data sample, corresponding to an integrated luminosity of $1.0 \text{ fb}^{-1}$, collected with the LHCb detector, the following relative branching fractions are measured

$$\frac{B(B^+ \to p\bar{p}K^+)}{B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 4.91 \pm 0.19 \text{ (stat) } \pm 0.14 \text{ (syst)},$$

$$\frac{B(B^+ \to p\bar{p}K^+)}{B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 2.02 \pm 0.10 \text{ (stat) } \pm 0.08 \text{ (syst)},$$

$$\frac{B(B^+ \to \eta_c(1S)K^+ \to p\bar{p}K^+)}{B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 0.578 \pm 0.035 \text{ (stat) } \pm 0.025 \text{ (syst)},$$

$$\frac{B(B^+ \to \eta_c(2S)K^+ \to p\bar{p}K^+)}{B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} = 0.080 \pm 0.012 \text{ (stat) } \pm 0.009 \text{ (syst)}.$$

An upper limit on the ratio $\frac{B(B^+ \to X(3872)K^+ \to p\bar{p}K^+)}{B(B^+ \to J/\psi K^+ \to p\bar{p}K^+)} < 0.017$ is obtained, from which a limit of

$$\frac{B(X(3872) \to p\bar{p})}{B(X(3872) \to J/\psi\pi^+\pi^-)} < 2.0 \times 10^{-3}$$

is derived.

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