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Bounds for spectral projectors on the Euclidean cylinder

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Abstract. We prove essentially optimal bounds for norms of spectral projectors on thin spherical shells for the Laplacian on the cylinder $(\mathbb{R}/\mathbb{Z}) \times \mathbb{R}$. In contrast to previous investigations into spectral projectors on tori, having one unbounded dimension available permits a compact self-contained proof.

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1. Introduction

1.1. Spectral projectors on general manifolds and tori

Given a Riemannian manifold with Laplace–Beltrami operator Δ , consider the spectral projector $P_{\lambda,\delta}$ on (perhaps generalized) eigenfunctions with eigenvalues within $O(\delta)$ of λ . It is defined through functional calculus by the formula

$$P_{\lambda,\delta} = P_{\lambda,\delta}^{\chi} = \chi\left(\frac{\sqrt{-\Delta}-\lambda}{\delta}\right),$$

where χ is a cutoff function, which is irrelevant for our purposes.

An interesting question is to determine the operator norm from L^2 to L^p , with p > 2, of this operator. A theorem of Sogge [5] gives an optimal answer for any Riemannian manifold if $\delta = 1$

$$||P_{\lambda,1}||_{L^2 \to L^p} \lesssim \lambda^{\frac{d-1}{2} - \frac{d}{p}} + \lambda^{\frac{d-1}{2} \left(\frac{1}{2} - \frac{1}{p}\right)}.$$

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While this completely answers the question if $\delta > 1$, the case $\delta < 1$ is still widely open. Understanding the case $\delta < 1$ requires a global analysis on the Riemannian manifold, which makes it very delicate.

In the case of the rational torus $\mathbb{R}^d/\mathbb{Z}^d$, L^p bounds on eigenfunctions attracted a lot of attention; this corresponds to the choice $\delta = 1/\lambda$, since the distance between two consecutive eigenvalues is $\sim \frac{1}{\lambda}$. The best result in this direction is due to Bourgain and Demeter [3]. More recently, the authors of the present paper [4] considered the problem for general values of λ and δ , conjectured the bound for general tori

$$\left\|P_{\lambda,\delta}\right\|_{L^2 \to L^p} \lesssim \lambda^{\frac{d-1}{2} - \frac{d}{p}} \delta^{1/2} + (\lambda\delta)^{\frac{(d-1)}{2} \left(\frac{1}{2} - \frac{1}{p}\right)} \quad \text{for } \delta > 1/\lambda,$$

and were able to establish this bound for a range of the parameters δ , λ , p.

A full proof of this conjecture seems very challenging in every dimension *d*. Restricting to the case d = 2, consider the case $(\mathbb{R}/\mathbb{Z}) \times \mathbb{R} = \mathbb{T} \times \mathbb{R}$ instead of \mathbb{T}^2 . The conjecture remains identical, but a short proof, relying on ℓ^2 decoupling, can be provided; this is the main observation of the present paper. Generalizations to higher dimensions are certainly possible.

1.2. The Euclidean cylinder

On $\mathbb{T} \times \mathbb{R} = (\mathbb{R}/\mathbb{Z}) \times \mathbb{R}$, we choose coordinates (x, y), with $x \in [0, 1]$ and $y \in \mathbb{R}$. The Laplacian operator is given by

$$\Delta = \partial_x^2 + \partial_y^2$$

A function *f* on $\mathbb{T} \times \mathbb{R}$ can be expanded through Fourier series in *x* and Fourier transform in *y*:

$$f(x) = \sum_{k \in \mathbb{Z}} \int_{\mathbb{R}} \widehat{f}(k, \eta) e^{2\pi i (kx + \eta y)} d\eta$$

The spectral projector can then be expressed as

$$P_{\lambda,\delta}f(x) = \sum_{k \in \mathbb{Z}} \int_{\mathbb{R}} \chi\left(\frac{\sqrt{k^2 + \eta^2} - \lambda}{\delta}\right) \widehat{f}(k,\eta) e^{2\pi i \left(kx + \eta y\right)} d\eta.$$

Theorem 1 (Main theorem). *If* $\lambda > 1$ *and* $\delta < 1$ *,*

$$\|P_{\lambda,\delta}\|_{L^2 \to L^p} \lesssim_{\epsilon} \lambda^{\epsilon} \delta^{-\epsilon} \left[\lambda^{\frac{1}{2} - \frac{2}{p}} \delta^{\frac{1}{2}} + (\lambda\delta)^{\frac{1}{4} - \frac{1}{2p}}\right]$$

Furthermore, this estimate is optimal, up to the subpolynomial factor $\lambda^{\epsilon} \delta^{-\epsilon}$ and the multiplicative constant.

1.3. Strichartz estimates

It is interesting to draw a parallel with Strichartz estimates in dimension 2 for the Schrödinger equation, in which case the critical exponent equals 4. It was proved in the foundational paper of Bourgain [2] that

$$\left\|e^{it\Delta}f\right\|_{L^4\left([0,1]\times\mathbb{T}^2\right)}\lesssim_s \|f\|_{H^s\left(\mathbb{T}^2\right)} \quad \text{for } s>0.$$

Takaoka and Tzvetkov [6] proved that the above inequality fails for s = 0, but that, on $\mathbb{T} \times \mathbb{R}$,

$$\left\| e^{it\Delta} f \right\|_{L^4([0,1] \times \mathbb{T} \times \mathbb{R})} \lesssim \left\| f \right\|_{L^2(\mathbb{T} \times \mathbb{R})}.$$

Finally, Barron, Christ and Pausader [1] determined the correct global (in time) estimate, for which a further summation index is needed. These examples suggest that optimal estimates might differ by subpolynomial factors between \mathbb{T}^2 and $\mathbb{T} \times \mathbb{R}$.

2. Proof of the Theorem 1 (main theorem)

Proof. By Plancherel's theorem, it suffices to prove

$$\|f\|_{L^p} \lesssim_{\epsilon} \lambda^{\epsilon} \delta^{-\epsilon} \left[\lambda^{\frac{1}{2} - \frac{2}{p}} \delta^{\frac{1}{2}} + (\lambda \delta)^{\frac{1}{4} - \frac{1}{2p}}\right] \|f\|_{L^2}$$
(1)

for *f* a function whose Fourier transform is localized in the corona $\mathscr{C}_{\lambda,\delta}$ of radius λ and with thickness δ/λ :

$$\mathscr{C}_{\lambda,\delta} = \left\{ (k,\eta) \text{ such that } \lambda - \delta < \sqrt{k^2 + \eta^2} < \lambda + \delta \right\}.$$

By symmetry, one can furthermore assume that $\hat{f}(k,\eta)$ is localized in the first quadrant $k,\eta \ge 0$.

The function f can be split into two pieces, which will correspond to the two terms on the right-hand side of (1).

$$f(x) = \left| \sum_{|k-\lambda| \le \frac{1}{\delta}} + \sum_{|k-\lambda| > \frac{1}{\delta}} \right| \int_{\mathbb{R}} \widehat{f}(k,\eta) e^{2\pi i (kx+\eta y)} d\eta = f_1(x) + f_2(x).$$

The Case $|k - \lambda| \le \frac{1}{\delta}$. The Fourier support of f_1 is made up of a collection of segments. We will see in Lemma 2 below that the added length of these segments can be bounded by

$$|\operatorname{Supp} \widehat{f}_1| \lesssim \sqrt{\lambda \delta}$$

Therefore, by the Cauchy-Schwartz inequality,

$$\|f_1\|_{L^{\infty}} \leq |\operatorname{Supp} \widehat{f}_1|^{1/2} \cdot \|\widehat{f}_1\|_{L^2}^{1/2} \lesssim (\lambda \delta)^{1/4} \|f_1\|_{L^2}.$$

Interpolating with L^2 , this gives

$$\|f_1\|_{L^p} \lesssim (\lambda \delta)^{\frac{1}{4} - \frac{1}{2p}} \|f\|_{L^2}$$

The Case $|k - \lambda| > \frac{1}{\delta}$. We start by choosing a function $\phi \in \mathscr{S}$ which is > 1/2 on [-1,1], and has Fourier support in [-1,1]. We use periodicity in the *x* variable to expand the range of *x* from $x \in [0,1]$ to $x < \delta^{-1}$, so that

$$\left\|f_2\right\|_{L^p(\mathbb{T}\times\mathbb{R})} \lesssim \delta^{1/p} \left\|\phi\left(\delta x\right)\sum_{|k-\lambda|>\frac{1}{\delta}} \int_{\mathbb{R}} \widehat{f}(k,\eta) e^{2\pi i (kx+\eta y)} d\eta\right\|_{L^p(\mathbb{R}^2)}$$

We now change variables as follows: $X = \lambda x$, $Y = \lambda y$, $K = k/\lambda$, $H = \eta/\lambda$,

$$f_2(x,y) = \lambda F(X,Y), \qquad F(X,Y) = \phi\left(\frac{\delta X}{\lambda}\right) \sum_{\substack{K \in \mathbb{Z}/\lambda \\ |K-1| > \frac{1}{\delta \lambda}}} \int_{\mathbb{R}} \widehat{f}(\lambda K, \lambda H) e^{2\pi i (KX+HY)} \, dH$$

to obtain

$$\|f_2\|_{L^p(\mathbb{T}\times\mathbb{R})} \lesssim \delta^{1/p} \lambda^{1-\frac{2}{p}} \|F\|_{L^p(\mathbb{R}^2)}.$$

The effect of this change of variables is that the function of (X, Y) whose L^p norm we want to evaluate has Fourier transform supported in the corona $\mathcal{C}_{1,3\delta/\lambda}$ of radius 1 and width δ/λ , and also in the first quadrant $X, Y \ge 0$. This enables us to apply the ℓ^2 decoupling theorem of Bourgain and Demeter [3]: for a smooth partition of unity (χ_{θ}) corresponding to a suitable almost disjoint covering of $\mathcal{C}_{1,3\delta/\lambda}$ by caps (θ) of size $\sim \frac{\delta}{\lambda} \times \sqrt{\frac{\delta}{\lambda}}$,

$$\|f_2\|_{L^p(\mathbb{T}\times\mathbb{R})} \lesssim \delta^{\frac{1}{p}} \lambda^{1-\frac{2}{p}} \left(\delta/\lambda\right)^{-\frac{1}{4}+\frac{3}{2p}-\epsilon} \left(\sum_{\theta} \|\chi_{\theta}(D)F\|_{L^p(\mathbb{R}^2)}^2\right)^{1/2},$$

where $\chi_{\theta}(D)$ is the Fourier multiplier with symbol χ_{θ} .

We now apply the inequality $\|g\|_{L^p(\mathbb{R}^d)} \lesssim \|g\|_{L^2} |\operatorname{Supp} \widehat{g}|^{\frac{1}{2} - \frac{1}{p}}$ (if $p \ge 2$), which follows by applying in turn the Hausdorff–Young and Hölder inequalities, and then the Plancherel theorem. Since by Lemma 2 below

$$\left|\operatorname{Supp} \widehat{\chi_{\theta}(D)F}\right| = \left|\operatorname{Supp} \chi_{\theta}\widehat{F}\right| \lesssim \delta^{5/2} \lambda^{-3/2},$$

it follows that

$$\|f_2\|_{L^p(\mathbb{T}\times\mathbb{R})} \lesssim \delta^{\frac{1}{p}} \lambda^{1-\frac{2}{p}} (\delta/\lambda)^{-\frac{1}{4}+\frac{3}{2p}-\epsilon} (\delta^{5/2}\lambda^{-3/2})^{\frac{1}{2}-\frac{1}{p}} \left(\sum_{\theta} \|\chi_{\theta}(D)F\|_{L^2(\mathbb{R}^2)}^2\right)^{1/2}.$$

By almost orthogonality, this becomes

$$\|f_2\|_{L^p(\mathbb{T}\times\mathbb{R})} \lesssim \delta^{\frac{1}{p}} \lambda^{1-\frac{2}{p}} (\delta/\lambda)^{-\frac{1}{4}+\frac{3}{2p}-\epsilon} (\delta^{5/2} \lambda^{-3/2})^{\frac{1}{2}-\frac{1}{p}} \|F\|_{L^2(\mathbb{R}^2)}$$

Finally, undoing the change of variables and using once again periodicity in the x variable gives

$$\|f_2\|_{L^p(\mathbb{T}\times\mathbb{R})} \lesssim \delta^{\frac{1}{p}} \lambda^{1-\frac{2}{p}} (\delta/\lambda)^{-\frac{1}{4}+\frac{3}{2p}-\epsilon} (\delta^{5/2}\lambda^{-3/2})^{\frac{1}{2}-\frac{1}{p}} \delta^{-1/2} \|f\|_{L^2}$$

$$\lesssim \lambda^{\frac{1}{2}-\frac{2}{p}} \delta^{1/2}$$

Optimality. The optimality of the statement of the theorem is proved through two examples. The first one is an analog of the Knapp example: assume $\lambda \in \mathbb{N}$, and consider the function *g* given by its Fourier transform

$$\widehat{g}(k,\eta) = \mathbf{1}_{\lambda}(k)\chi\left(\frac{\eta}{\sqrt{\lambda\delta}}\right).$$

Here, $\mathbf{1}_{\lambda}$ is the indicator function of $\{\lambda\}$ and χ is a cutoff function with a sufficiently small support, so that Supp $\widehat{g} \subset \mathscr{C}_{\lambda,\delta}$. In physical space,

$$g(x, y) = \sqrt{\lambda \delta} e^{2\pi i \lambda x} \widehat{\chi} \left(\sqrt{\lambda \delta} y \right).$$

It has L^p norm ~ $(\lambda\delta)^{\frac{1}{2}-\frac{1}{2p}}$, so that

$$\frac{\|g\|_{L^p}}{\|g\|_{L^2}} \sim (\lambda \delta)^{\frac{1}{4} - \frac{1}{2p}}.$$

We now consider the function h given by its Fourier transform

$$\widehat{h}(k,\eta) = \mathbf{1}_{\mathscr{C}_{\lambda,\delta}}(k,\eta)\mathbf{1}_{[0,\lambda/2]}(k);$$

here, $\mathbf{1}_{\mathscr{C}_{\lambda,\delta}}$ is the indicator function of the annulus, and $\mathbf{1}_{[0,\lambda/2]}$ the indicator function of the interval. It is easy to check that $|\operatorname{Supp} \hat{h}| \sim \lambda \delta$, so that $\|h\|_{L^{\infty}} \sim \lambda \delta$ and $\|h\|_{L^{2}} \sim \sqrt{\lambda \delta}$, and finally

$$\frac{\|h\|_{L^{\infty}}}{\|h\|_{L^2}} \sim \sqrt{\lambda \delta}.$$

By the Bernstein inequality,

$$\frac{\|h\|_{L^p}}{\|h\|_{L^2}} \gtrsim \lambda^{-2/p} \frac{\|h\|_{L^\infty}}{\|h\|_{L^2}} \sim \lambda^{\frac{1}{2} - \frac{2}{p}} \delta^{1/2}.$$

The examples *g* and *h* show that the statement of the Theorem 1 is optimal, up to subpolynomial losses. \Box

3. Bounds for the Fourier support

Lemma 2 (Bound on the size of Fourier support). With the notations of the proof of Theorem 1,

(i) The function f_1 is a function on $\mathbb{T} \times \mathbb{R}$. As such, its Fourier transform is supported on a union of lines, and has one-dimensional measure

$$|\operatorname{Supp} \widehat{f}_1| \lesssim \sqrt{\lambda \delta}.$$

(ii) The function $\chi_{\theta}(D)F$ is a function on \mathbb{R}^2 . Its Fourier transform is defined on \mathbb{R}^2 , and has two-dimensional measure

$$|\operatorname{Supp} \chi_{\theta} \widehat{F}| \lesssim \delta^{5/2} \lambda^{-3/2}.$$

Proof.

(i) Consider *f* as in the proof of Theorem 1, namely with Fourier support in $\mathscr{C}_{\lambda,\delta}$. Since (k,η) range in $\mathbb{Z} \times \mathbb{R}$ with $k, \eta \ge 0$, the Fourier support of *f* is contained in $\cup_{k \in \mathbb{Z}} \{k\} \times E_k^{\lambda}$, where

$$E_k^{\lambda} = \begin{cases} \emptyset & (k \ge k_+), \\ \left(0, \sqrt{k_+^2 - k^2}\right) & (|k - \lambda| < \delta), \\ \left(\sqrt{(\lambda - \delta)^2 - k^2}, \sqrt{k_+^2 - k^2}\right) & (0 \le k \le \lambda - \delta) \end{cases}$$

Recalling that \hat{f}_1 is just \hat{f} restricted to $|k - \lambda| \le \frac{1}{\delta}$, one can then add up these pieces to get the bound

$$\left|\operatorname{Supp} \widehat{f}_{1}\right| \lesssim \sum_{\max\{0, \lambda - \frac{1}{\delta}\} \leq k < k_{+}} |E_{k}| \lesssim \sqrt{\lambda \delta} + \sum_{\substack{k \geq 0 \\ \frac{1}{\delta} \geq \lambda - k > \delta}} \frac{\delta \lambda}{\sqrt{\lambda(\lambda - x)}}$$

and as $y \mapsto 1/\sqrt{y}$ is decreasing this is

$$\leq 2\sqrt{\lambda\delta} + \int_{\delta}^{\min\{\lambda, \frac{1}{\delta}\}} \frac{\delta\lambda}{\sqrt{\lambda y}} \, dy \leq 4\sqrt{\lambda\delta}.$$

(ii) Turning to F, it has Fourier support in

$$\bigcup_{\substack{k \in \mathbb{Z} \\ k-\lambda| > \frac{1}{\delta}}} \left[\frac{k}{\lambda} - \frac{2\delta}{\lambda}, \frac{k}{\lambda} + \frac{2\delta}{\lambda} \right] \times D_k^{\lambda}, \qquad D_k^{\lambda} = \left\{ H, 1 - \frac{\delta}{\lambda} < \sqrt{\frac{k^2}{\lambda^2} + H^2} < 1 + \frac{\delta}{\lambda} \right\}.$$
(2)

Consider $\chi_{\theta}(D)F$, for a cap θ with dimensions

$$\sim \frac{\delta}{\lambda} \times \sqrt{\frac{\delta}{\lambda}}$$

adapted to the corona $\mathscr{C}_{1,3\delta/\lambda}$. Given such a cap, there is $j \in \mathbb{N}$ with $2^j > \frac{1}{\delta}$ such that every point in the intersection of θ with the set (2) satisfies $|\lambda - k| \sim 2^j$.

There are around $\sqrt{\delta}2^{j/2}$ such values of k for which the vertical strip $\left[\frac{k}{\lambda} - \frac{2\delta}{\lambda}, \frac{k}{\lambda} + \frac{2\delta}{\lambda}\right]$ intersects the cap θ . For each such k, the size of D_k^{λ} is $\sim \delta \lambda^{-1/2} 2^{-j/2}$. Hence, adding up the contributions in (2),

$$\left|\operatorname{Supp} \chi_{\theta}(D)F\right| \lesssim \sqrt{\delta} 2^{j/2} \cdot \delta \lambda^{-1} \cdot \delta \lambda^{-1/2} 2^{-j/2} = \delta^{5/2} \lambda^{-3/2}.$$

References

- [1] A. Barron, M. Christ, B. Pausader, "Global endpoint strichartz estimates for Schrödinger equations on the cylinder $\mathbb{R} \times \mathbb{T}$ ", Nonlinear Anal., Theory Methods Appl. **206** (2021), article no. 112172.
- [2] J. Bourgain, "Fourier transform restriction phenomena for certain lattice subsets and applications to nonlinear evolution equations. II: The KdV-equation", *Geom. Funct. Anal.* 3 (1993), no. 3, p. 209-262.
 [3] J. Bourgain, C. Demeter, "The proof of the l² decoupling conjecture", *Ann. Math.* 182 (2015), no. 1, p. 351-389.
- [4] P. Germain, S. L. R. Myerson, "Bounds for spectral projectors on tori", https://arxiv.org/abs/2104.13274v1, 2021.
- [5] C. D. Sogge, Fourier integrals in classical analysis, Cambridge Tracts in Mathematic, vol. 105, Cambridge University Press, 1993.
- [6] H. Takaoka, N. Tzvetkov, "On 2D nonlinear schrödinger equations with data on $\mathbb{R} \times \mathbb{T}$ ", J. Funct. Anal. 182 (2001), no. 2, p. 427-442.