## Supporting Information

## Octanuclear Heterobimetallic $\{Ni_4Ln_4\}$ Assemblies Possessing $Ln_4$ Square Grid $[2\times2]$ Motifs: Synthesis, Structure and Magnetism

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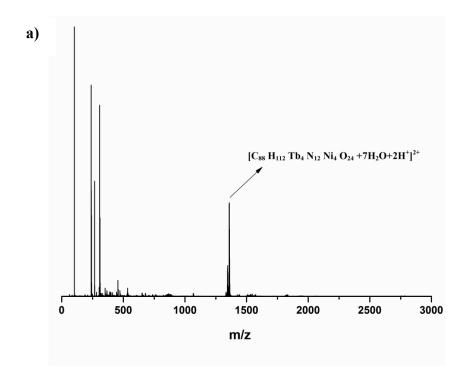
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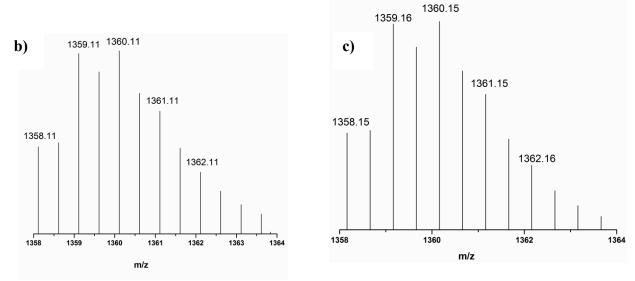
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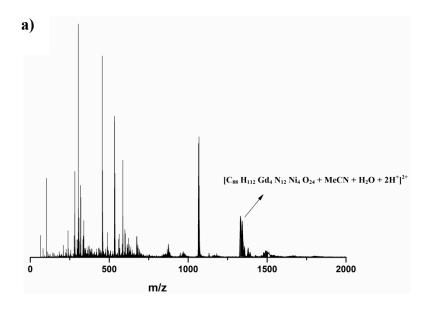
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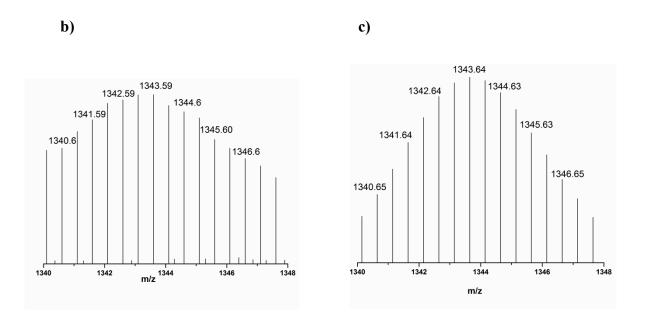
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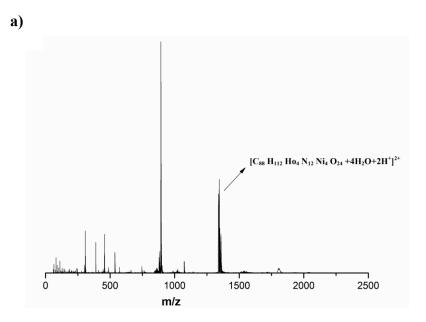


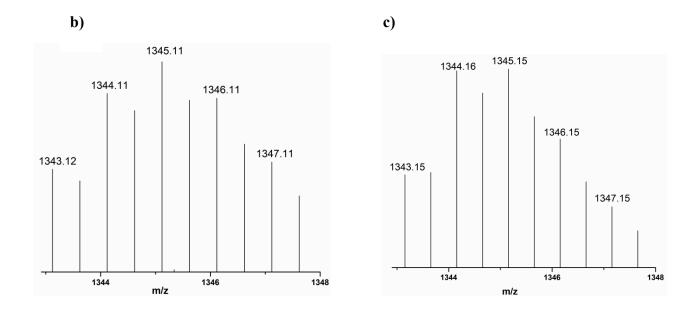
**FigureS1**. a) Full range ESI-MS spectrum of complex 2. b) Experimental and c) Simulated mass spectral pattern of the species  $[C_{88} H_{116}Tb_4N_{12}Ni_4O_{24} + 7H_2O - 2H^+]^{2+}$ .



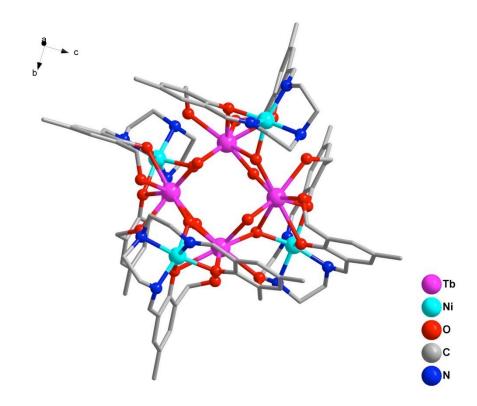


**Figure S2**. **a)** Full range ESI-MS spectrum of complex **3**. **b)** Experimental and **c)** Simulated mass spectral pattern of the species  $[C_{88} H_{116}Gd_4N_{12}Ni_4O_{24} + MeCN + H_2O - 2H^+]^{2+}$ .





**Figure S3.a)** Full range ESI-MS spectrum of complex **4**. **b)** Experimental and **c)** Simulated mass spectral pattern of the species  $[C_{88} H_{116} Ho_4 N_{12} Ni_4 O_{24} + 4H_2 O - 2H^+]^{2+}$ .

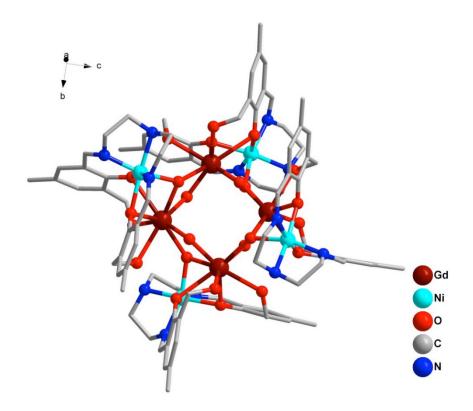


**Figure S4**. Molecular Structure of compound **2** (hydrogen atoms, chlorides and non coordinated solvent molecules were omitted for clarity).

TableS1. Selected Bond Distances (Å) and Bond angles (°) for compound2

Coordination environ around	Selected bond lengths	Selected bond angles
N1 03 06 06 05	Ni(1)-N(3) = 2.028(8) Ni(1)-N(1) = 2.046(8) Ni(1)-O(3) = 2.086(7) Ni(1)-O(5) = 2.100(7) Ni(1)-O(6) = 2.115(7) Ni(1)-N(2) = 2.164(9)	Ni(1)-O(5)-Tb(1)* = 139.2(3) Ni(1)-O(3)-Tb(1) = 92.3(2) Ni(1)-OH(6)-Tb(1) = 84.8(2)

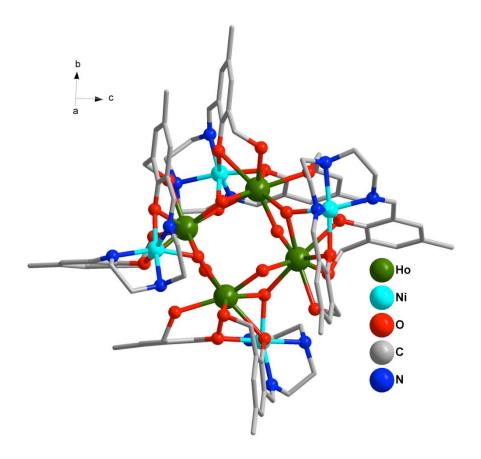
	Tb(1)-O(2)* = 2.244(6)	Tb(1)*-O(8)-Tb(1)=112.9(3)
O2	Tb(1)-O(2) = 2.269(6)	Tb(1)-O(6)-Tb(1)*=101.8(2)
	Tb(1)-O(3) = 2.374(6)	
06	Tb(1)-O(5) = 2.390(7)	
	Tb(1)-O(4) = 2.411(7)	
05	Tb(1)-O(5)*=2.440(7)	
O5* Tb1	Tb(1)-O(1) = 2.469(7)	
	Tb(1)-O(6) = 2.631(6)	
	Tb(1)-Ni(1) = 3.223(13)	
03		



**Figure S5**. Molecular Structure of compound **3** (selected hydrogen atoms, chlorides and non-coordinated solvent molecules were omitted for clarity).

TableS2. Selected Bond Distances (Å) and Bond angles (°) for compound 3.

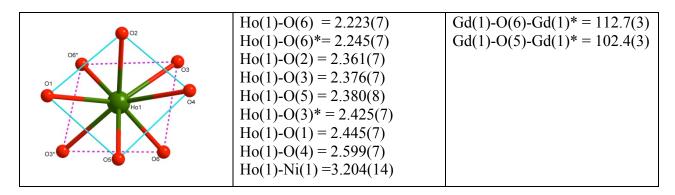
Coordination environ around	Selected bond lengths	Selected bond angles
N1 N3 N3 O3	Ni(1)-N(3) = 2.014(9) $Ni(1)-N(2) = 2.047(9)$ $Ni(1)-O(3) = 2.090(7)$ $Ni(1)-O(5) = 2.094(7)$ $Ni(1)-O(2) = 2.115(7)$ $Ni(1)-N(1) = 2.161(9)$	Ni(1)-OH(5)-Gd(1) = 91.8(3) Ni(1)-OH(5)-Gd(1)* = 139.4(4) Ni(1)-O(3)-Gd(1) = 91.7(3)
O5 O5 O5 O5 O5	Gd(1)-O(6) = 2.250(6) $Gd(1)-O(6)* = 2.290(7)$ $Gd(1)-O(5) = 2.396(7)$ $Gd(1)-O(3) = 2.405(7)$ $Gd(1)-O(1) = 2.443(7)$ $Gd(1)-O(5)* = 2.455(7)$ $Gd(1)-O(4) = 2.501(7)$ $Gd(1)-O(2) = 2.645(7)$ $Gd(1)-Ni(1) = 3.232(14)$	Gd(1)-O(6)-Gd(1)* = 112.7(3) Gd(1)-OH(5)-Gd(1)* = 102.4(3)

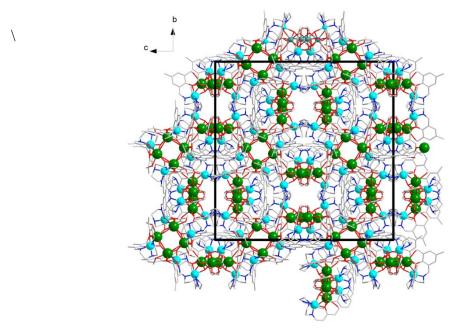


**Figure S6**. Molecular Structure of compound **4** (selected hydrogen atoms, chlorides and non coordinated solvent molecules were omitted for clarity).

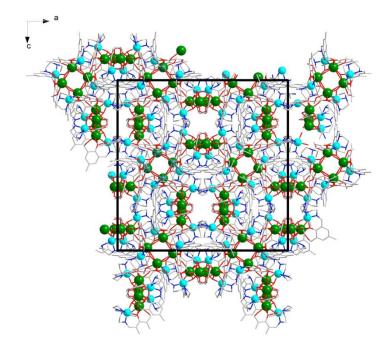
TableS3. Selected Bond Distances (Å) and Bond angles (°) for compound 4.

O2 O3 Ni 1 Ni	Ni(1)-N(2) = 2.040(9) Ni(1)-N(3) = 2.050(9) Ni(1)-O(2) = 2.078(7) Ni(1)-O(4) = 2.106(7) Ni(1)-O(3) = 2.109(7) Ni(1)-N(1) = 2.164(9)	Ni(1)-O(3)-Ho(1) = 91.0(3) Ni(1)-O(3)-Ho(1)* = 138.5(3) Ni(1)-O(2)-Ho(1) = 92.2(3)
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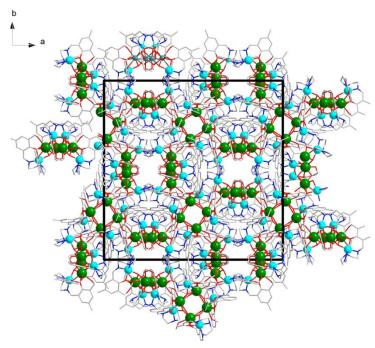




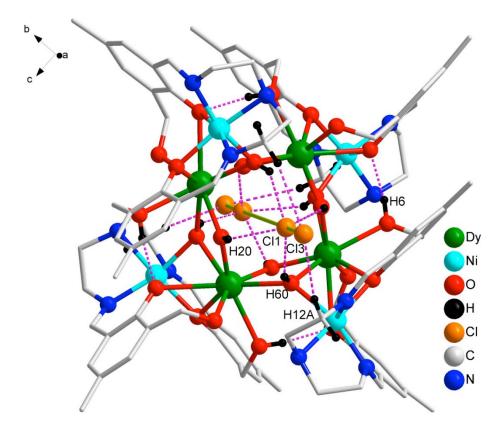
**Figure S7.** Crystal packing along the x axis.



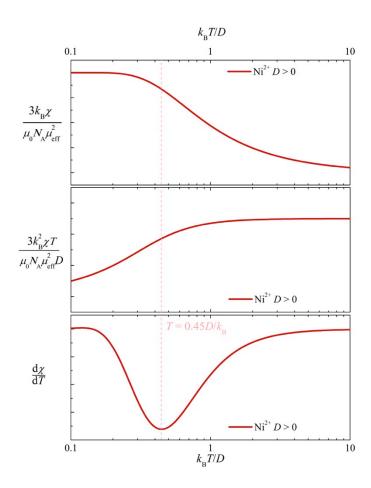
**Figure S8.** Crystal packing along the y axis.



**Figure S9.** Crystal packing along the z axis.



**Figure S10**. Molecular Structure of compound **1** (selected hydrogen atoms and non-coordinated solvent molecules were omitted for clarity).



**Figure S11**. Simulated dc magnetic susceptibility curves of normalized  $\chi$ ,  $\chi T$  and  $d\chi/dT$  against normalized temperature for Ni<sup>2+</sup> ions of easy-plane single-ion anisotropy. A derived relation between the  $d\chi/dT$  minimum and Ni<sup>2+</sup> single-ion anisotropy,  $D_{\text{Ni}}$ , is indicated.

## **Magnetic properties**

The susceptibility ( $\chi$ ) of an ensemble of spin-1 Ni<sup>2+</sup> ions in an applied field (H) was considered using the Hamiltonian

$$\mathcal{H} = D_{\text{Ni}} S_z^2 + g \mu_B \mu_0 \mathbf{S} \cdot \mathbf{H}, \quad (S1)$$

where  $D_{\text{Ni}}$  is the single-ion anisotropy, S is the Ni<sup>2+</sup> spin and g is the (isotropic) g-factor.

For a powdered sample, the susceptibility can be deduced from an exact diagonalization of Eq.  $S1^{1}$  for fields applied parallel and perpendicular to the magnetic hard-axis (z), yielding

$$\chi = \frac{N_{\rm A}\mu_0\mu_{\rm B}^2g^2S(S+1)}{3}\left(\frac{e^{-\beta D}}{k_BT(1+2e^{-\beta D})} + \frac{2(1-e^{-\beta D})}{D(1+2e^{-\beta D})}\right),$$

where  $\beta = 1/k_BT$  and S = 1 for Ni<sup>2+</sup> in an octahedral coordination environment. The resultant susceptibility for  $D_{\text{Ni}} > 0$  (Figure S11) indicates that a significant thermal depopulation of crystal-field split Ni<sup>2+</sup> levels causes a decrease in the product  $\chi T$  on cooling concurrent with a pronounced minimum in  $d\chi/dT$  centred at  $T = 0.45 D_{\text{Ni}}/k_B$ .

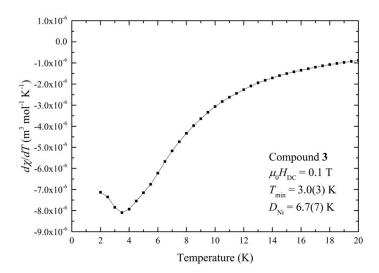
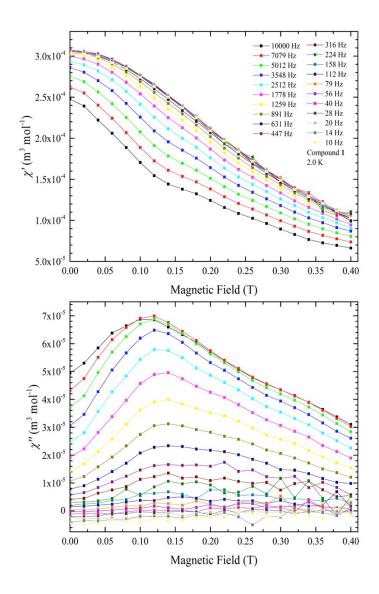
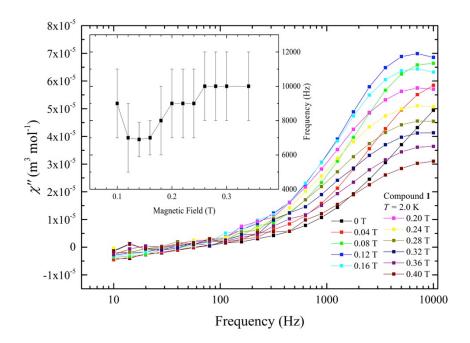


Figure S12. Dc magnetic susceptibility of 3 (Ni<sub>4</sub>Gd<sub>4</sub>) differentiated with respect to T,  $d\chi/dT$ , after interpolating data in 0.5 K steps and smoothing the resultant data via 3 point adjacent averaging. The minimum temperature is noted as well as the resultant  $D_{Ni}$  value via the relation from Figure S11.



**Figure S13**. Magnetic field study of  $Ni_4Dy_4$  at 2 K via ac susceptibility measurements. A dc magnetic field was increased from 0 T to 0.4 T in 0.02 T steps with a full set of frequencies measured at each step.



**Figure S14**. Out-of-phase susceptibility data from Figure S13 of Ni<sub>4</sub>Dy<sub>4</sub> at 2 K plotted against frequency. Data at all magnetic fields approach a maximum at high frequencies. Inset: Frequency position of maxima  $\chi$ '' plotted as a function of magnetic field.

When slow relaxation of magnetization is present in a system, the complex magnetic susceptibility may be described by the generalized Debye model:

$$\chi = \chi_S + \frac{\chi_T + \chi_S}{1 + [i\omega\tau(T)]^{1-\alpha}}, \quad (S2)$$

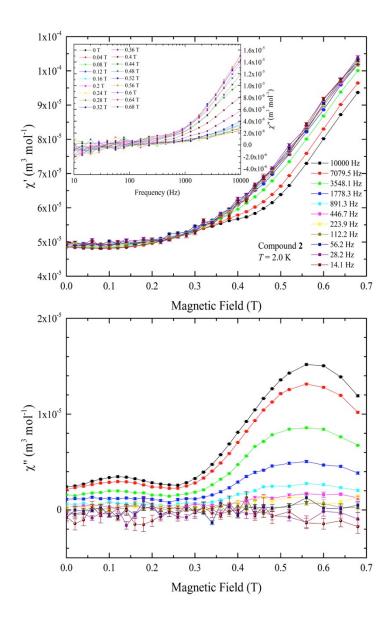
where  $\chi_S$  and  $\chi_T$  are the adiabatic and isothermal susceptibility respectively, and  $0 \le \alpha \le 1$  measures the spread of the relaxation time,  $\tau$ .<sup>32</sup> Rearranging the above to eliminate  $\omega \tau$  results in<sup>2</sup>

$$\chi'' = -\frac{(\chi_{\rm T} - \chi_{\rm S})}{2 \tan \left[\frac{\pi}{2} (1 - \alpha)\right]} + \sqrt{(\chi' - \chi_{\rm S})(\chi_{\rm T} - \chi') + \frac{(\chi_{\rm T} - \chi_{\rm S})^2}{4 \left(\tan \left[\frac{\pi}{2} (1 - \alpha)\right]\right)^2}}.$$
 (S3)

This slow relaxation may be detected in ac susceptibility data as a series of peaks in  $\chi''$  at the points at which the condition  $\omega \tau = 1$  is satisfied.<sup>2</sup> For an organometallic cluster system  $\tau$  may be expected to follow

$$\frac{1}{\tau} = aTB^2 + \frac{b}{1 + cB^2} + dT^n + \tau_0 e^{\frac{-U_{\text{eff}}}{kT}}, \quad (S4)$$

where a, b, c, d and  $\tau_0$  are constants and B is the applied magnetic field.<sup>3</sup> Each term represents a different possible relaxation process: direct, quantum tunnelling, Raman and Orbach processes respectively.<sup>3</sup> SMM behavior is described by the Orbach term, which dominates other, none magnetic-field dependent terms in SMMs.<sup>4</sup> To best detect the Orbach process, ac susceptibility measurements can be performed at the optimum dc magnetic field that minimizes the field-dependent terms. This is shown in the inset to Figure S14 in which the frequency position of the maxima in  $\chi''$  have been plotted against dc magnetic field. A minimum is seen at approximately  $\mu_0 H_{\rm DC} = 0.14$  T, which corresponds to the optimum field for the observation of SMM behavior in Ni<sub>4</sub>Dy<sub>4</sub>.



**Figure S15**. Magnetic field study of **2** (Ni<sub>4</sub>Tb<sub>4</sub>) at 2 K. A dc magnetic field was increased from 0 T to 0.68 T in initially 0.02 T steps before switching to 0.04 T steps above 0.48 T. Twenty-one frequencies were measured at each field. Only select frequencies are shown for clarity. Inset: out-of-phase susceptibility of Ni<sub>4</sub>Tb<sub>4</sub> at 2 K plotted versus frequency. No maxima are visible within the accessible frequency range.

The field study of Ni<sub>4</sub>Tb<sub>4</sub> (Figure S15) shows a number of differences compared to Ni<sub>4</sub>Dy<sub>4</sub>.  $\chi'$  first decreases slightly on application of a magnetic field up to 0.1 T, before increasing more rapidly and displaying a marked frequency separation. A series of maxima are seen in  $\chi''$  in the low-field region at approximately the same fields at which  $\chi'$  exhibits a minimum, with the higher frequency peaks occurring at slightly higher fields, which is in contrast to the maxima seen in the Ni<sub>4</sub>Dy<sub>4</sub> material. As the dc magnetic field is increased,  $\chi''$  forms another, larger set of peaks. The  $\chi''$  data of Ni<sub>4</sub>Tb<sub>4</sub> is plotted as constant field curves against frequency in the inset to Figure S15. This plot is very similar to the corresponding Ni<sub>4</sub>Dy<sub>4</sub> plot (Figure S13). However, if maxima are to be observed in these data they lie outside the measured frequency range.

- 1) Molecular Magnetism, Kahn, O; Wiley-VCH, New York, 1993.
- 2) Hagiwara, M. J. Magn. Magn. Mater. 1998, 177, 89-90.
- 3) Novikov, V. V.; Pavlov, A. A.; Nelyubina, Y. V.; Boulon, M.-E.; Varzatskii, O. A.; Voloshin, Y. Z.; Winpenny, R. E. P. *J. Am. Chem. Soc.* **2015**, *137*, 9792-9795.
- 4) Gatteschi, D.; Sessoli, R. Angew. Chem. Int. Ed. 2003, 42, 268-297.